

Assignment 6: MCMC Simulation and Review Problems

Due Monday July 30 at 11:59PM in Instructor’s Inbox

Preface: In this assignment, you will use computational software to implement the Markov Chain Monte Carlo (MCMC) simulation algorithm we simulated ourselves in class.

1. MCMC of Spins in a Magnetic Field

In this problem, we simulate spin system depicted in Fig. 1.

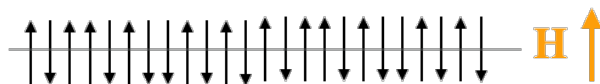


Figure 1: Particular microstate of the system we are simulating.

Using the partition function, we can show that for a collection of N spins each of which has magnetic dipole μ and subject to magnetic field H , the average spin is given by

$$\langle s \rangle = \tanh \left(\frac{\mu H}{k_B T} \right), \quad (1)$$

where $\tanh(x) = (e^x - e^{-x}) / (e^x + e^{-x})$. We will use a simulation to compare the theoretical result Eq.(1) to one obtains from computation.

We will be using *Mathematica* to complete this problem, so we need to establish some code preliminaries.

- (i) Log in to your account in one of MIT’s Athena Clusters, and go to the MITES 2018 –Physics III course website.
- (ii) Download the code `spin_model_simulation.nb` from the course webpage and open it in *Mathematica*.
- (iii) Select a block of code and run it by pressing `Shift+Enter`.

Now we begin the problem

- (a) Run each line of the “Function Definitions” section of the code to determine what the associated function does
- (b) **Low Temperature Simulation**

Fill in all the “. . .” in the code. In particular define/fill in the following:

- Parameter definition $\mu = 1.0$, $H = 1.0$, $k_B T = 0.1$
- There are $N_{\text{spins}} = 1000$ in the system
- There are $N_{\text{steps}} = 10,000$
- An initial spin configuration with N_{spins} all of which have value -1 ;

- Energy of microstate is given by

$$E(\{s_i\}) = -\mu H \sum_{i=1}^N s_i \quad (2)$$

- The ratio between Boltzmann factors for final and initial spin configurations

$$\exp \left[-\frac{1}{k_B T} (E(\{s_i\}_{\text{final}}) - E(\{s_i\}_{\text{initial}})) \right] \quad (3)$$

- The theoretical value of $\langle s \rangle$ obtained from Eq.(1)

(c) **Low Temperature Simulation**

Perform a new simulation (i.e., all the parts of (b)) at the higher temperature $k_B T = 10.0$. *Hint: Copy and paste is your friend here.*

- (d) In the simulations, why do we only take the final 100 time steps, to compute the average spin?
- (e) In (b) and (c), how do the theoretical values of $\langle s \rangle$ compare with the computational results? How might we improve the correspondence between the two.

Submitting: As your submission for this part of the assignment, you should first make sure you write your name at the top of the notebook. Save the notebook as `lastname_spin_model_2018.nb` and email it to `mwilliams@physics.harvard.edu`.

2. **Optional Bonus: MCMC of Mean-Field Ising Model**

Simulate the Mean-Field Ising Model for $J = 1.0$ and *computationally* compute average spin at the two temperature $k_B T = 0.5$ and $k_B T = 1.5$. Compare these computational results with Figure 12(b) in Lecture Notes 05.

Review Problems

(Not to submit)

Preface: This assignment contains a list of review problems to help you prepare for the final exam.

1. Protein Expression and Probability

For a particular model of a gene in a cell, the probability density that said gene produces a concentration of x proteins during the cell cycle is given by

$$p(x) = A \left(\frac{x}{b}\right)^N e^{-x/b}, \quad (4)$$

where b is a biological constant with units of concentration and A is a normalization parameter.

- The concentration of proteins that can be produced ranges from zero to infinite. What must A be in order for Eq.(4) to be normalized?
- What is the mean of the normalized probability density?
- What is the standard deviation of the normalized probability density?

2. Adsorption

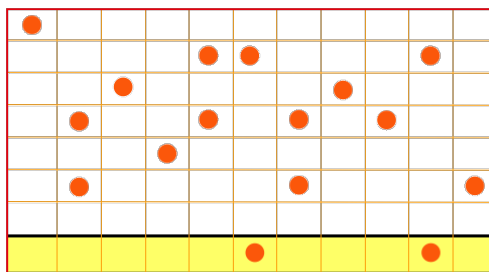


Figure 2: Particular microstate of system. The possible surface states are shown in yellow. The possible bulk states are the remaining states.

A gas is in a container, and some of the atoms are stuck (i.e., “adsorped”) to the surface of the container, and others are free to move in the volume. We will study this system using a discrete model of the positions the atoms can take on.

Let M be the number of positions an atom can occupy while in the bulk, and let N be the number of positions an atom can occupy while on the surface. We have N atoms so that the surface states can be completely filled. Let $k \leq N$ be the number of atoms that are on the surface. A *single* atom has energy $-\varepsilon$ while it is on the surface and energy 0 while it is in the bulk. M , N , and ε are all constants, and k is a random variable. M , N , and k are all much greater than 1.

- How many microstates does the system have for a particular value of k ?
- What is the free energy of this system? *You can write this free energy in terms of factorials.*
- Use Stirling’s approximation $\ln N! \simeq N \ln N - N$ to approximate the result in (b). Derive **one** condition that k must satisfy in order to define the thermal equilibrium of the system. *We often state two conditions associated with thermal equilibrium. Only apply the first-derivative condition.*

3. Nucleotide zipper

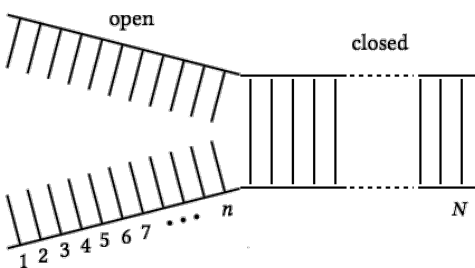


Figure 3: "Zipper" of nucleotides

A strand of DNA can be modeled as a zipper with links between N base-pairs. Each link has the energy 0 in the closed state and the energy ε the open state. Assume the zipper can only open from one end, so that the n th link can open only if all the links 1, 2, ..., $n - 1$ are open. Assume that the system is at a temperature T .

- How many microstates does this system have? What are the energies of these microstates?
- Find the partition function for a single DNA strand as a function of T and ε . *Hint: Evaluate the summation using the geometric series identity $\sum_{n=0}^N x^n = (1 - x^{N+1})(1 - x)^{-1}$.*
- Assume $N \gg 1$. Find the average number of open links as a function of T and ε .

4. Entropy of an Ideal Gas

For a system in thermal equilibrium defined by the partition function Z , we determined that the Helmholtz free energy of the system is

$$F = -k_B T \ln Z = -k_B T \ln \sum_{\{i\}} e^{-\beta E_i}, \quad (5)$$

where i denotes a microstate of the system, and E_i is the energy of that microstate.

- For a system in thermal equilibrium, the Gibbs entropy is given by

$$S = -k_B T \sum_{\{i\}} p_i \ln p_i = -k_B \sum_{\{i\}} \frac{e^{-\beta E_i}}{Z} \ln \left(\frac{e^{-\beta E_i}}{Z} \right). \quad (6)$$

Given Eq.(5), show that

$$S = -\frac{\partial F}{\partial T}. \quad (7)$$

Note: $\beta = 1/k_B T$.

- Using Eq.(7), show that the entropy of an ideal gas for $N \gg 1$ is

$$S \simeq N k_B \left[\ln \left(\frac{V}{N} \right) - \frac{3}{2} \ln \left(\frac{h^2}{2\pi m k_B T} \right) + \frac{5}{2} \right]. \quad (8)$$

Note: You will need to use Stirling's approximation $\ln N! \simeq N \ln N - N$

5. **Partition Function as an Integral**

The partition function for a system of N spins is given by

$$Z(T) = \int_{-\infty}^{\infty} d\rho \exp [-\beta N f(\rho, T)], \tag{9}$$

where

$$f(\rho, T) = -a(T - T_c)\rho^2 + c\rho^4. \tag{10}$$

The parameters a , c , and T_c are constants in the model. T stands in for temperature.

Approximate the partition function using Laplace's method. For what values of T is the approximation invalid?

6. **Transition probabilities and equilibrium probabilities**

A person is in one of four rooms, and after each time step has a probability $1/(M + 1)$ of moving to an adjacent room where M is the number of adjacent rooms; The probability of remaining in place is $1/(M + 1)$.

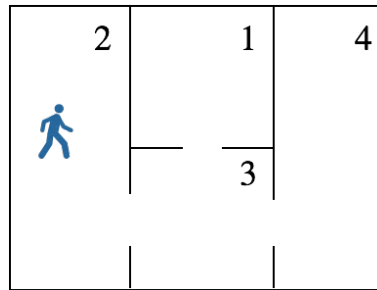


Figure 4: Four connected rooms. The person has a probability $1/(M + 1)$ of remaining in place and a probability of $1/(M + 1)$ of moving to any one of the adjacent rooms.

(a) Letting $\pi_{i \rightarrow j}$ represent the probability of transitioning from room i to room j in a single time step, fill in the elements below

$$\begin{array}{cccc}
 \pi_{1 \rightarrow 1} = & \pi_{1 \rightarrow 2} = & \pi_{1 \rightarrow 3} = & \pi_{1 \rightarrow 4} = \\
 \pi_{2 \rightarrow 1} = & \pi_{2 \rightarrow 2} = & \pi_{2 \rightarrow 3} = & \pi_{2 \rightarrow 4} = \\
 \pi_{3 \rightarrow 1} = & \pi_{3 \rightarrow 2} = & \pi_{3 \rightarrow 3} = & \pi_{3 \rightarrow 4} = \\
 \pi_{4 \rightarrow 1} = & \pi_{4 \rightarrow 2} = & \pi_{4 \rightarrow 3} = & \pi_{4 \rightarrow 4} =
 \end{array} \tag{11}$$

(b) Using the relationship between equilibrium probabilities and transition probabilities (along with the fact that the equilibrium probabilities are normalized), compute the equilibrium probabilities p_1^{eq} , p_2^{eq} , p_3^{eq} , and p_4^{eq} , to be in rooms 1, 2, 3, and 4.