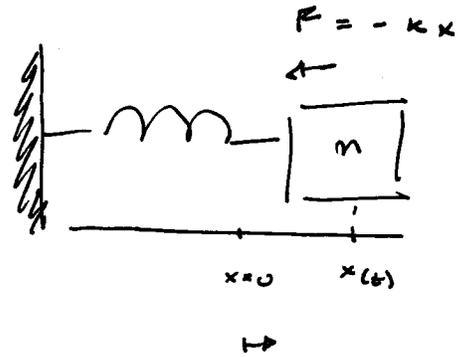


# AGENDA

June 28, 2010  
MITES 2010  
PHYSICS III

- [1] RECAP
- [2] LINEARITY & GEN SOLNS
- [3] COMPLEX SOLUTION & COMPLEX #S
- [4] COUPLED OSCILLATIONS



## [1] RECAP

WE WERE TRYING TO SOLVE THE SHM EQUATION

$$\ddot{x} + \omega^2 x = 0$$

WE GUESSED

$$x_1(t) = \sin(\alpha t) \quad \text{and FOUND}$$

$$0 = \ddot{x}_1(t) + \omega^2 x_1(t) = -\alpha^2 \sin(\alpha t) + \omega^2 \sin(\alpha t) \\ = (-\alpha^2 + \omega^2) \sin(\alpha t)$$

SO  $\alpha = \pm \omega$      $x_1(t) = \pm \sin(\alpha t)$

WE ALSO REALIZED WE COULD USE  $x_2(t) = \cos(\alpha t)$  BECAUSE

$$\dot{x}_2(t) = -\alpha \sin(\alpha t) \\ \ddot{x}_2(t) = -\alpha^2 \cos(\alpha t)$$

$$\ddot{x}_2 + \omega^2 x_2 = -\alpha^2 \cos(\alpha t) + \omega^2 \cos(\alpha t) \\ = (-\alpha^2 + \omega^2) \cos(\alpha t)$$

SO  $\alpha = \pm \omega$      $x_2(t) = \cos(\omega t)$

WE FOUND TWO SOLUTIONS. HOW DO WE FIND MORE? (1)

## [2] LINEARITY & GEN. SOLNS

OUR TWO SOLNS ARE  $x_1(t)$  &  $x_2(t)$  WHICH SOLVE THE DIFFERENTIAL EQUATION AND THEREFORE SATISFY

$$\ddot{x}_1(t) + \omega^2 x_1(t) = 0$$

$$\ddot{x}_2(t) + \omega^2 x_2(t) = 0$$

HOW CAN WE GET OTHER SOLUTIONS?

1) ADD!

LET  $x_3(t) = x_1(t) + x_2(t)$

THEN

$$\begin{aligned}\ddot{x}_3(t) + \omega^2 x_3(t) &= \ddot{x}_1(t) + \ddot{x}_2(t) + \omega^2 (x_1(t) + x_2(t)) \\ &= \ddot{x}_1(t) + \omega^2 x_1(t) + \ddot{x}_2(t) + \omega^2 x_2(t) \\ &= 0 + 0 = 0\end{aligned}$$

SO  $x_3(t)$  IS ANOTHER SOLUTION

WHAT ELSE CAN WE DO?

2) MULTIPLY BY A CONSTANT!

LET  $x_4(t) = A x_1(t)$       $x_5(t) = B x_2(t)$

THEN

$$\ddot{x}_4(t) + \omega^2 x_4(t) = A \ddot{x}_1 + \omega^2 A x_1$$

$$= A (\ddot{x}_1 + \omega^2 x_1)$$

$$= 0$$

(Same as  $x_1(t)$ )

SO WE CAN ADD TWO SOLUTIONS OR MULTIPLY THEM BY CONSTANTS TO OBTAIN OTHER SOLUTIONS. (2)

THE MOST GENERAL SOLUTION IS A COMBINATION OF BOTH OPERATIONS

GENERAL SOLUTION

IF  $x_1(t)$  &  $x_2(t)$  ARE LINEARLY INDEPENDENT SOLUTIONS OF A D.E. THEN THE MOST GENERAL SOLUTION IS

$$x(t) = C_1 x_1(t) + C_2 x_2(t)$$

OR

$$= \boxed{A x_1(t) + B x_2(t)}$$

WHERE  $A$  &  $B$  ARE ARBITRARY CONSTANTS DIFFERENTLY DEFINED BY A SPECIFIC SYSTEM.

SO FOR SHO

$$x_1(t) = \cos(\omega t)$$

$$x_2(t) = \sin(\omega t)$$

THE MOST GENERAL SOLUTION IS

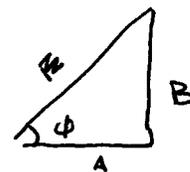
$$\boxed{x(t) = A \cos(\omega t) + B \sin(\omega t)}$$

HOW ELSE CAN WE WRITE THIS?

- SWITCH THE INDEPENDENT ~~CONSTANT~~ CONSTANTS  $A$  &  $B$  FOR TWO OTHER CONSTANTS  $E$  &  $\phi$ , WHERE

$$A = E \cos \phi$$

$$B = E \sin \phi$$



THEN

$$\begin{aligned} x(t) &= E \cos \phi \cos(\omega t) + E \sin \phi \sin(\omega t) \\ &= E (\cos \phi \cos(\omega t) + \sin \phi \sin(\omega t)) \end{aligned}$$

FROM TRIG, WE KNOW

$$\cos(\alpha + \beta) = \cos \alpha \cos \beta - \sin \alpha \sin \beta$$

$$\cos(\alpha - \beta) = \cos \alpha \cos \beta + \sin \alpha \sin \beta$$

SO

$$x(t) = E \cos(\omega t - \phi)$$

NOTICE; WE STILL HAVE TWO INDEPENDENT CONSTANTS  $E$  &  $\phi$

THIS IS A GENERAL RESULT.

NTH ORDER DIFFERENTIAL EQUATION



N ARBITRARY & INDEPENDENT CONSTANTS

BUT!

THERE IS STILL ANOTHER WAY TO WRITE  $x(t)$

[3] COMPLEX & COMPLEX SOLUTION

INSTEAD OF A TRIG FUNCTION LET'S GUESS AN EXPONENTIAL

NAMELY

$$z(t) = e^{\alpha t}$$

\* WE USE THE LABEL "z" BECAUSE  $z(t)$  WILL TURN OUT TO BE A COMPLEX #.

PLUG THIS INTO OUR S.O.E.

$$\begin{aligned} \ddot{z}(t) + \omega^2 z(t) &= \alpha^2 e^{\alpha t} + \omega^2 e^{\alpha t} \\ &= (\alpha^2 + \omega^2) e^{\alpha t} = 0 \end{aligned}$$

$$\text{SO } \alpha^2 = -\omega^2$$

$$\alpha = \pm \sqrt{-\omega^2}$$

$$= \pm \sqrt{-1} \cdot \sqrt{\omega^2} = \pm i\omega \quad \text{WHERE}$$

$$i = \sqrt{-1}$$

SO WE HAVE TWO SOLUTIONS  $e^{+i\omega t}$ ,  $e^{-i\omega t}$ . IT TURNS OUT THAT THESE TWO ARE LINEARLY INDEPENDENT.

SO THE GENERAL SOLUTION IS

$$z(t) = C e^{i\omega t} + D e^{-i\omega t}$$

HOW IS THIS RESULT EQUIVALENT TO OUR FIRST?

ANS. THROUGH EULER'S FORMULA

$$e^{i\theta} = \cos \theta + i \sin \theta$$

HOW IS THIS TRUE?

- USUALLY PROVED WITH TAYLOR SERIES

- WE WILL TAKE IT AS FACT

\* SUGGESTION: SET  $f(\theta) = \cos \theta + i \sin \theta$ . DIFFERENTIATE TO FIND  $\frac{df}{d\theta} = i f$ . SOLVE DIFFERENTIAL EQUATION.

SO OUR  $z(t)$  CAN BE WRITTEN AS.

$$z(t) = C e^{i\omega t} + D e^{-i\omega t}$$

$$= C (\cos(\omega t) + i \sin(\omega t)) + D (\cos(\omega t) - i \sin(\omega t))$$

$$= \underbrace{(C + D)}_{\text{Complex } \#} \cos(\omega t) + i \underbrace{(C - D)}_{\text{Complex } \#} \sin(\omega t)$$

Complex #

Complex #

↓

↓

$$z(t) = (A_1 + iA_2) \cos(\omega t) + (B_1 + iB_2) \sin(\omega t)$$

$$\text{Re}(z) = A_1 \cos(\omega t) + B_1 \sin(\omega t) = x(t)$$

CONCLUSION:  $x(t)$  IS CONTAINED IN THE REAL COMPONENT (5)  
OF  $z(t)$ . SIDENOTE: IT IS ALSO CONTAINED IN THE IMAGINARY COMPONENT

SO WE HAVE THREE FOUR WAYS OF WRITING THE SOLUTION TO

$$\ddot{x} + \omega^2 x = 0$$

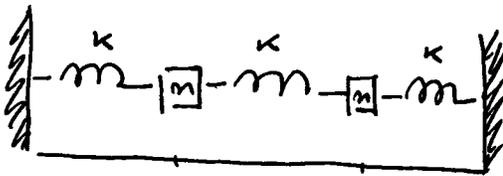
⇔

$$x(t) = C e^{i\omega t} + D e^{-i\omega t} \quad \text{COMPLEX}$$

$$\text{Real Part.} = A \cos(\omega t) + B \sin(\omega t) \quad \text{REAL}$$

$$= R \cos(\omega t - \phi) \quad \text{REAL}$$

### [4] COUPLED OSCILLATIONS



TWO MASSES CONNECTED TO EACH OTHER WITH A SPRING OF CONSTANT  $k$  AND TO AN ADJACENT WALL WITH A SPRING OF CONSTANT  $k$ .

- WHAT IS THE RESULTING MOTION?

↳ REST LENGTH IS AT

$$x_1 = x_2 = 0$$

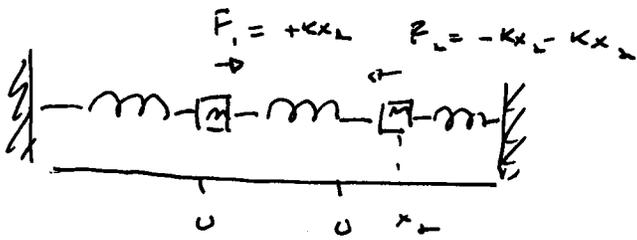
### PROCEDURE

- 1) DERIVE EOMS.
- 2) SOLVE.

WE DERIVE THE EOM'S FROM NEWTON'S SECOND LAW AND A CONSIDERATION OF FORCES.

CONSIDERING THE MOTION OF BOTH MASSES IS DIFFICULT. LET'S CONSIDER ONE AT A TIME

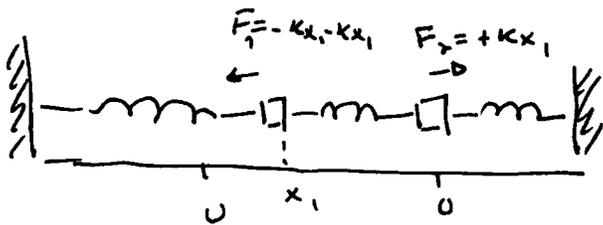
LET  $x_1(t) = 0$  ; LET  $x_2(t) > 0$



$$F_1 = +kx_2$$

$$F_2 = -2kx_2$$

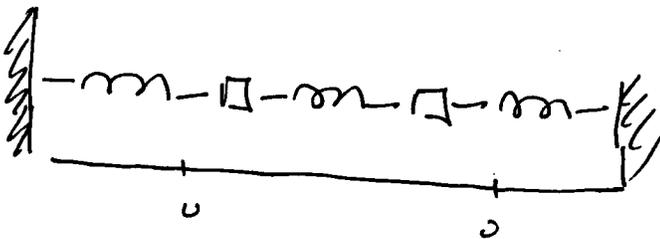
LET  $x_2(t) = 0$  ; LET  $x_1(t) > 0$



$$F_1 = -2kx_1$$

$$F_2 = kx_1$$

GENERAL CASE : LET  $x_1(t) \neq x_2(t) \neq 0$  (SUM OF BOTH CASES)



$$F_1 = -2kx_1 + kx_2$$

$$F_2 = kx_1 - 2kx_2$$

NEWTON'S SECOND LAW

$$F_1 = m\ddot{x}_1 = -2kx_1 + kx_2$$

$$F_2 = m\ddot{x}_2 = kx_1 - 2kx_2$$

LET  $\frac{k}{m} = \omega^2$

$$\begin{cases} \ddot{x}_1 = -2\omega^2 x_1 + \omega^2 x_2 \\ \ddot{x}_2 = \omega^2 x_1 - 2\omega^2 x_2 \end{cases}$$

EOMS

WE HAVE OUR EOMS, NOW WE "ONLY" NEED TO SOLVE FOR  $x_1(t)$  &  $x_2(t)$

WE WILL DO THIS IN TWO WAYS

1) MATRIX ~~METHOD~~ METHOD

2) LINEAR COMBINATION METHOD.

1) MATRIX METHOD

FROM MATRIX MULTIPLICATION

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} e \\ f \end{pmatrix} \Leftrightarrow \begin{cases} ax + by = e \\ cx + dy = f \end{cases}$$

SO

$$\ddot{x}_1 = -2\omega^2 x_1 + \omega^2 x_2$$

$$\ddot{x}_2 = \omega^2 x_1 - 2\omega^2 x_2$$

$$\begin{pmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{pmatrix} = \begin{pmatrix} -2\omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

LIKE BEFORE, WE WILL GUESS AN EXPONENTIAL SOLUTION

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} A \\ B \end{pmatrix} e^{\alpha t}$$

$$\frac{d^2}{dt^2}$$

$$\begin{pmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{pmatrix} = \alpha^2 \begin{pmatrix} A \\ B \end{pmatrix} e^{\alpha t}$$

SO THE EOM NOW READS

$$\begin{pmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{pmatrix} = \alpha^2 \begin{pmatrix} A \\ B \end{pmatrix} e^{\alpha t} = \begin{pmatrix} -2\omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} e^{\alpha t}$$

$$\alpha^2 \begin{pmatrix} A \\ B \end{pmatrix} e^{\alpha t} = \begin{pmatrix} -2\omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} e^{\alpha t}$$

EXPONENTIALS CANCEL & WE MULTIPLY THE LEFT SIDE BY THE IDENTITY MATRIX TO GET BOTH SIDES INTO THE "2x2 . 2x1" MATRIX MATRIX FORMAT.

$$\alpha^2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = \begin{pmatrix} -2\omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

MUVE EVERYTHING TO ONE SIDE

$$0 = \begin{pmatrix} -2\omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} - \alpha^2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

$$= \left[ \begin{pmatrix} -2\omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 \end{pmatrix} - \begin{pmatrix} \alpha^2 & 0 \\ 0 & \alpha^2 \end{pmatrix} \right] \begin{pmatrix} A \\ B \end{pmatrix}$$

$$= \begin{pmatrix} -2\omega^2 - \alpha^2 & \omega^2 \\ \omega^2 & -2\omega^2 - \alpha^2 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix}$$

$$= \hat{M} \vec{a}$$

so  $\boxed{\hat{M} \vec{a} = \vec{0}}$

THIS WAS A LONG PROCESS SO LET'S REMIND OURSELVES OF WHAT WE'RE LOOKING FOR.

WE GUESSED  $\begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} A \\ B \end{pmatrix} e^{\alpha t}$  WE WANT TO KNOW  $A, B, \alpha, \kappa$

SO OUR EQUATION IS

$$\hat{M} \vec{a} = 0$$

HOW CAN WE SOLVE THIS? ANS: ASSUME WE KNOW  $\hat{M}^{-1}$  WHERE  $\hat{M}^{-1}$  IS THE INVERSE OF  $\hat{M}$  AND THEREFORE SATISFIES

$$\hat{M}^{-1} \hat{M} = \mathbb{1} \leftarrow \text{Identity MATRIX}$$

THEN, WE MULTIPLY BOTH SIDES BY ~~THE~~  $\hat{M}^{-1}$  OF  $\hat{M} \vec{a} = 0$  BY  $\hat{M}^{-1}$  AND WE FIND

$$\hat{M}^{-1} \hat{M} \vec{a} = \hat{M}^{-1} 0$$

$$\mathbb{1} \vec{a} = 0$$

$$\vec{a} = 0 = \begin{pmatrix} A \\ B \end{pmatrix}$$

SO WE GET  $A = B = 0$  AND NO MOTION EXISTS. THIS IS CLEARLY NOT WHAT WE'RE LOOKING FOR.

SO WE DO NOT WANT  $\hat{M}^{-1}$  TO EXIST.

WE CAN ENFORCE THIS BY MAKING  $\det \hat{M} = 0$ .

IT TURNS OUT

$$\hat{M}^{-1} = \frac{1}{\det \hat{M}} \left( \text{"SOME MATRIX"} \right)$$

SO IF  $\det \hat{M} = 0$  THEN  $\hat{M}^{-1}$  DOES NOT EXIST AND WE OBTAIN A NONZERO SOLUTION. SO WE WANT

$$\boxed{\det \hat{M} \neq 0}$$

WE HAVE A  $2 \times 2$  MATRIX SO WE WILL NEED THE FORMULA FOR THE DETERMINANT OF A  $2 \times 2$  MATRIX

WE HAVE THE  $\pm$  CASE SO ONE SOLUTION IS FOUND TO BE

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(1)} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} \left( A_1 e^{+2\omega t} + A_2 e^{-2\omega t} \right)$$

Now for  $x_2$

$$\alpha_2 = \pm i\sqrt{3}\omega$$

$$\hat{M}_2 = \begin{pmatrix} -2\omega^2 - \alpha^2 & \omega^2 \\ \omega^2 & -2\omega^2 - \alpha^2 \end{pmatrix} = \begin{pmatrix} -2\omega^2 + 3\omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 + 3\omega^2 \end{pmatrix} = \begin{pmatrix} \omega^2 & \omega^2 \\ \omega^2 & \omega^2 \end{pmatrix}$$

so

$$\hat{M}_2 \vec{x} \Rightarrow \begin{pmatrix} \omega^2 & \omega^2 \\ \omega^2 & \omega^2 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = \begin{aligned} \omega^2 A + \omega^2 B &= 0 \\ \omega^2 A + \omega^2 B &= 0 \end{aligned}$$

$$\boxed{A = -B}$$

so we find that our other solution is

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(2)} = \begin{pmatrix} 1 \\ -1 \end{pmatrix} \left( B_1 e^{+i\sqrt{3}\omega t} + B_2 e^{-i\sqrt{3}\omega t} \right)$$

Using our various ways of writing the solution to the SHO from we know we may write

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(1)} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} E_1 \cos(\omega t - \phi_1) \quad \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(2)} = \begin{pmatrix} 1 \\ -1 \end{pmatrix} E_2 \cos(\sqrt{3}\omega t - \phi_2)$$

IF

$$\hat{O} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

Then  $\det \hat{O} = ad - bc$

SO

$$\hat{M} = \begin{pmatrix} -2\omega^2 - a^2 & \omega^2 \\ \omega^2 & -2\omega^2 - a^2 \end{pmatrix} \Rightarrow \det \hat{M} = (-2\omega^2 - a^2)^2 - \omega^4 = 0$$

SO

$$(-2\omega^2 - a^2)^2 - \omega^4 = 0$$

$$-\omega^4 + 4\omega^4 + 4\omega^2 a^2 + a^4 = 0$$

$$a^4 + 3\omega^4 + 4\omega^2 a^2 = 0$$

$$(a^2 + \omega^2)(a^2 + 3\omega^2) = 0$$

↓

$$\boxed{a_1 = \pm i\omega \quad a_2 = \pm \sqrt{3}i\omega}$$

THESE DEFINE THE CHARACTERISTIC FREQUENCIES OF THE MOTION.

WHAT DO  $a_1$  &  $a_2$  SAY ABOUT  $A$  &  $B$ ?

LET'S CHOOSE  $a_1$  FIRST

$$\underline{a_1 = \pm i\omega}$$

$$\hat{M}_1 = \begin{pmatrix} -2\omega^2 - a_1^2 & \omega^2 \\ \omega^2 & -2\omega^2 - a_1^2 \end{pmatrix} = \begin{pmatrix} -2\omega^2 + \omega^2 & \omega^2 \\ \omega^2 & -2\omega^2 + \omega^2 \end{pmatrix} = \begin{pmatrix} -\omega^2 & \omega^2 \\ \omega^2 & -\omega^2 \end{pmatrix}$$

SO WITH

$$\hat{M}_1 \vec{r} = 0 \Rightarrow \begin{pmatrix} -\omega^2 & \omega^2 \\ \omega^2 & -\omega^2 \end{pmatrix} \begin{pmatrix} A \\ B \end{pmatrix} = 0$$

$$-\omega^2 A + \omega^2 B = 0$$

$$\omega^2 A - \omega^2 B = 0$$

$$\Rightarrow \boxed{A = B}$$

(12)  
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WHERE WHEN MOVING FROM COMPLEX EXPONENTIALS TO TRIG FUNCTIONS WE TOOK THE REAL PART OF THE EXPONENTIALS.

THESE TWO SOLUTIONS

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(1)} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} E_1 \cos(\omega t - \phi_1)$$

2.

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(2)} = \begin{pmatrix} 1 \\ -1 \end{pmatrix} E_2 \cos(\sqrt{3}\omega t - \phi_2)$$

ARE CALLED THE "NORMAL ~~COORDINATES~~ MODES" OF THE MOTION. THEY PHYSICALLY REPRESENT THE SIMPLEST NON-TRIVIAL MOTION OF THE SYSTEM. AND CONSEQUENTLY EVERY MOTION CAN BE REPRESENTED AS A LINEAR COMBINATION OF THESE NORMAL MODES.

THE MOST GENERAL SOLUTION IS A SUM OF THESE MODES

$$\begin{aligned} \begin{pmatrix} x_1(t) \\ x_2(t) \end{pmatrix} &= \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(1)} + \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}_{(2)} \\ &= \begin{pmatrix} 1 \\ 1 \end{pmatrix} E_1 \cos(\omega t - \phi_1) + \begin{pmatrix} 1 \\ -1 \end{pmatrix} E_2 \cos(\sqrt{3}\omega t - \phi_2) \end{aligned}$$

NOW FOR THE OTHER METHOD.

2) LINEAR COMBINATION

WE RETURN TO OUR EOMS

$$\ddot{x}_1 = -2\omega^2 x_1 + \omega^2 x_2$$

$$\ddot{x}_2 = \omega^2 x_1 - 2\omega^2 x_2$$

WE WILL FORM LINEAR COMBINATIONS (SUMS & DIFFERENCES)  
OF OUR TWO EQUATIONS, TO OBTAIN A SIMPLER RESULT.

1) SUM

$$\begin{aligned}\ddot{x}_1 + \ddot{x}_2 &= -2\omega^2 x_1 + \omega^2 x_2 + \omega^2 x_1 - 2\omega^2 x_2 \\ &= -\omega^2 x_1 - \omega^2 x_2 = -\omega^2 (x_1 + x_2)\end{aligned}$$

LET

$$u_+ = x_1 + x_2$$

SO WE HAVE

$$\ddot{u}_+ = -\omega^2 u_+$$

WHICH IS THE SHO EOM. FROM OUR PREVIOUS ANALYSIS  
WE KNOW THAT THE RESULT HAS THE SOLUTION.

$$u_+(t) = E_+ \cos(\omega t - \phi_+)$$

WHERE  $E_+$  &  $\phi_+$   
ARE CONSTANTS

2) DIFFERENCE

$$\begin{aligned}\ddot{x}_1 - \ddot{x}_2 &= -2\omega^2 (x_1 - x_2) + \omega^2 (x_2 - x_1) \\ &= -2\omega^2 (x_1 - x_2) - \omega^2 (x_1 - x_2) \\ &= -3\omega^2 (x_1 - x_2)\end{aligned}$$

$$\text{LET } u_- = x_1 - x_2$$

SO WE HAVE

$$\ddot{u}_- = -3\omega^2 u_-$$

WITH THE SOLUTION

$$u_-(t) = E_- \cos(\omega\sqrt{3} \omega t - \phi_-)$$

SO, WE FOUND  $u_+$  &  $u_-$ , BUT WHAT ARE  $x_1$  &  $x_2$ ?

$$u_+ = x_1 + x_2$$

$$u_- = x_1 - x_2$$

$$\Rightarrow \begin{cases} \frac{u_+ + u_-}{2} = x_1 \\ \frac{u_+ - u_-}{2} = x_2 \end{cases}$$

SO

$$x_1 = \frac{1}{2} u_+ + \frac{1}{2} u_-$$

$$= \frac{E_+}{2} \cos(\omega t - \phi_+) + \frac{E_-}{2} \cos(\omega\sqrt{3} \omega t - \phi_-)$$

$$x_2 = \frac{1}{2} u_+ - \frac{1}{2} u_-$$

$$= \frac{E_+}{2} \cos(\omega t - \phi_+) - \frac{E_-}{2} \cos(\omega\sqrt{3} \omega t - \phi_-)$$

WHICH CAN BE WRITTEN AS

$$\begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} \frac{E_+}{2} \cos(\omega t - \phi_+) + \begin{pmatrix} 1 \\ -1 \end{pmatrix} \frac{E_-}{2} \cos(\omega\sqrt{3} \omega t - \phi_-)$$

THIS RESULT IS EQUIVALENT TO THE ONE OBTAINED  
BY THE MATRIX METHOD.