

Favorable-Contact Model of Folding

Mobolaji Williams — Shakhnovich Group Meeting— Mar. 21, 2017

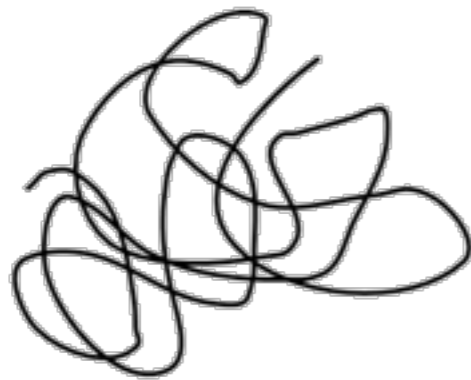
Connecting Protein Folding and Protein Design



Design and Folding Problems

How do we model the way structure determines sequence?

Study a model of possible sequence orders for a given presumed structure



Structure

Protein Design



Protein Folding

... A - R - H - G - L - H ...

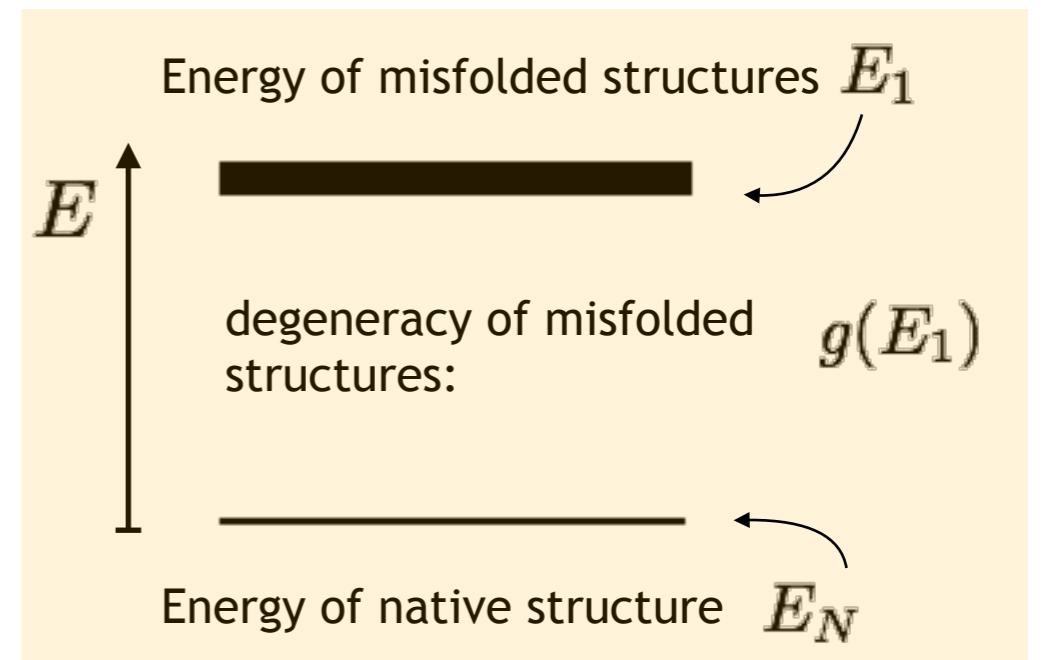
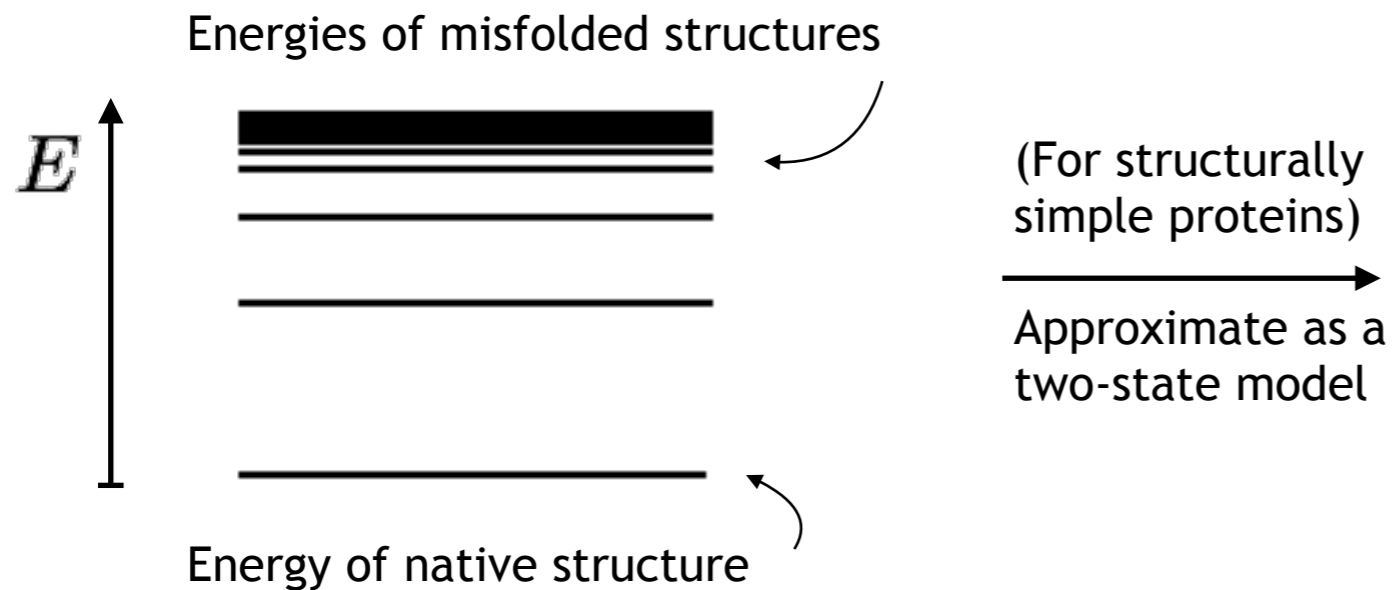
Sequence

How do we model the way sequence determines structure?

Study a model of possible pairwise contacts for a given sequence of contact regions

(Abstracted) Protein Design

Two-State Folding Model

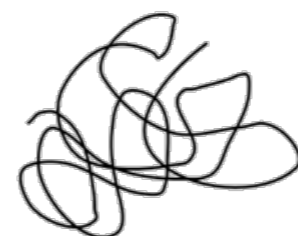


Stability of Native Structure

~ (Boltzmann Probability to be in Native Structure)

$$P_{\text{native}} = \frac{1}{1 + g(E_1)e^{-\beta(E_1 - E_N)}}$$

To find the sequence which maximizes energy gap, we search over sequence space for a given structure



Structure



- $\mathcal{S}_1 : R - D - \dots - L$
 - $\mathcal{S}_2 : D - R - \dots - L$
 - \vdots
 - $\mathcal{S}_{N_s} : V - Y - \dots - N$
- Possible AA Sequences

In order for sequence to properly fold the native structure must be highly stable (i.e., P_{native} must be as close to 1 as possible)

$E_1 - E_N \gg k_B T$: Energy gap between native and misfolded structures should be large

Protein Folding and Design

Problem with formalism:

*“For virtually any target structure, the lowest energy sequence will be a homopolymer, consisting of the amino acid with the largest self-attraction”**

**(Morrissey and Shakhnovich. "Design of proteins with selected thermal properties." (1996))*



Structure



Homopolymer?

This cannot be correct...

Suggested way forward:

Constrain amino acid composition

New Design Question:
Given a **fixed number of each type of amino-acid**
What sequence yields the lowest energy?

Implication of New Design Question

For a chain of length N , what is the size of our state space?



$$20^N \quad (\text{or } 2^N \text{ under a polar/nonpolar framing})$$



$$\frac{N!}{n_1!n_2!\dots n_{20}!} \quad (\text{or } \frac{N!}{n_{\text{polar}}!n_{\text{nonpolar}}!} \text{ under a polar/nonpolar framing})$$

We should search over the space of **permutations** of components

Partition Function of Permutations

Let's define the system more precisely and introduce a Hamiltonian.

System Definition:

- subunits are labeled as ω_i with $i = 1, \dots, N$
- the ordering of subunits with the zero energy is $\vec{\omega} \equiv (\omega_1, \omega_2, \dots, \omega_N)$
- an arbitrary state is $\vec{\theta}$ where $\vec{\theta} \in \{\text{perm}(\omega_1, \omega_2, \dots, \omega_N)\} \equiv \text{Sym}(\omega)$

Energy Definition:

- The state $\vec{\theta} = \vec{\omega}$ has zero energy and has subunits in the *correct* order. For all other states, there is an energy cost of λ_i for $\theta_i \neq \omega_i$.

$$\mathcal{H}_N(\{\theta_i\}) = \sum_{i=1}^N \lambda_i I_{\theta_i \neq \omega_i}$$

$$I_A = \begin{cases} 1 & A \text{ is true} \\ 0 & A \text{ is false} \end{cases}$$

Example: Three components

| State | Energy |
|----------------------------------|-------------------------------------|
| $(\omega_1, \omega_2, \omega_3)$ | 0 |
| $(\omega_2, \omega_1, \omega_3)$ | $\lambda_1 + \lambda_2$ |
| $(\omega_3, \omega_2, \omega_1)$ | $\lambda_1 + \lambda_3$ |
| $(\omega_1, \omega_3, \omega_2)$ | $\lambda_2 + \lambda_3$ |
| $(\omega_2, \omega_3, \omega_1)$ | $\lambda_1 + \lambda_2 + \lambda_3$ |
| $(\omega_3, \omega_1, \omega_2)$ | $\lambda_1 + \lambda_2 + \lambda_3$ |

Now let's compute the partition function

$$Z_N(\{\beta\lambda_i\}) = \sum_{\vec{\theta} \in \text{Sym}(\omega)} \exp\left(-\beta \sum_{i=1}^N \lambda_i I_{\theta_i \neq \omega_i}\right)$$

$$\rightarrow Z_N(\{\beta\lambda_i\}) = \int_0^\infty ds e^{-s} \prod_{\ell=1}^N [1 + (s-1)e^{-\beta\lambda_\ell}]$$

... so we can obtain a closed form expression ...

...but what does it mean?

Model of Permutations

Energy Definition:

- The state $\vec{\theta} = \vec{\omega}$ has zero energy and has subunits in the *correct* order. For all other states, there is an energy cost of λ_i for $\theta_i \neq \omega_i$.

$$\mathcal{H}_N(\{\theta_i\}) = \sum_{i=1}^N \lambda_i I_{\theta_i \neq \omega_i}$$

Partition Function:

$$Z_N(\{\beta\lambda_i\}) = \int_0^\infty ds e^{-s} \prod_{\ell=1}^N [1 + (s-1)e^{-\beta\lambda_\ell}]$$

What physics is contained in this partition function?

Progress comes from a simplification: $\lambda_i = \lambda_0$ for all i

“The same energy penalty for each subunit”

The partition function then simplifies to

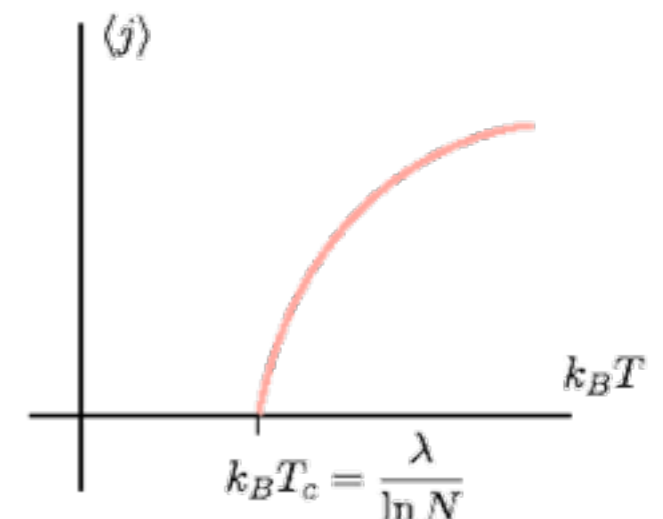
$$Z_N(\beta\lambda_0) = \int_0^\infty ds e^{-s} (1 + (s-1)e^{-\beta\lambda_0})^N$$

And the average number of incorrect components is

$$\langle j \rangle = \left\langle \sum_{i=1}^N I_{\theta_i \neq \omega_i} \right\rangle = \frac{\partial}{\partial(\beta\lambda_0)} \ln Z_N(\beta\lambda_0)$$

$$\longrightarrow \langle j \rangle \simeq \begin{cases} 0 & \text{for } T < T_c \\ N - e^{\beta\lambda_0} & \text{for } T > T_c \end{cases}$$

There is a transition temperature



→ Above a certain transition temperature, there is a spectrum of sequences (different from “correct sequence”) which yield the free energy minimum

(Abstracted) Protein Folding

Motivation: N-mer Problems Lattice Proteins

We can study protein-folding through **lattice protein models**

Properties of Lattice Protein Models

- **Geometry** is made irrelevant (all proteins have the same shape)
- **Stability** is determined via contact densities and interaction energy between lattice sites
- **Primary Structure:** Sequence of monomers defined by their hydrophobicity
- **Secondary Structure:** None
- **Tertiary Structure:** Configuration of chain in the space of the cube

27-monomer
lattice protein

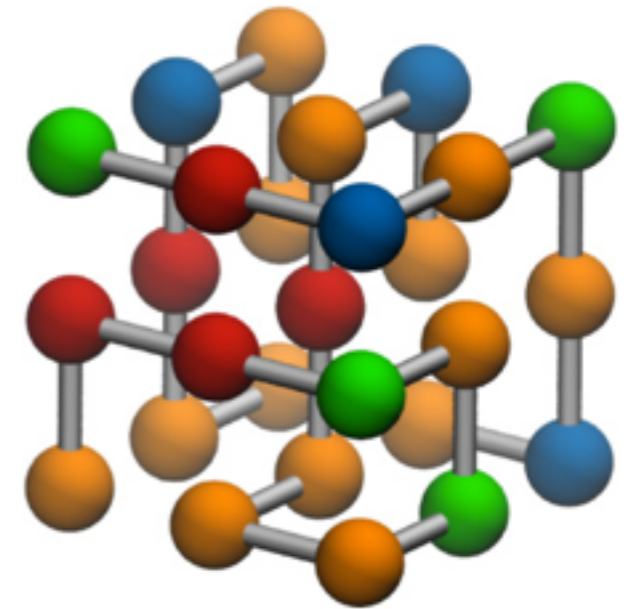


Figure from Whitford, Sanbonmatsu, and Onuchic.(2012)

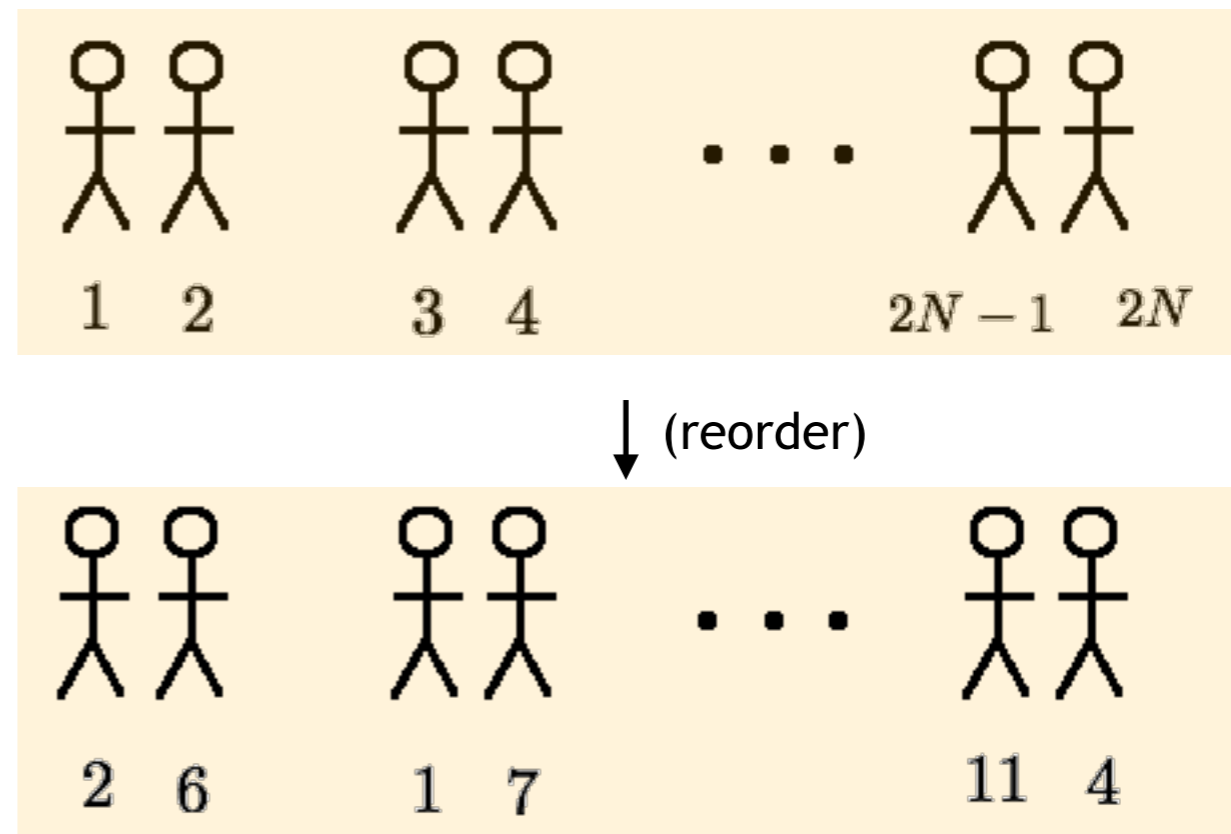
Configuration of chain
in the 3X3X3 space
represents a “folded”
structure

Color represents
hydrophobicity of
monomer

Can we consider other statistical
models of similar simplicity?

Dancing Partners Problem

- N partners (i.e., $2N$ total people) arrive at a dancing party
- Each dance partner pair is re-ordered such that each person may or may not be with their original dance partner.
- In how many ways can this happen? (What is the size of the state space?)



$$|\mathcal{S}_N| = \frac{(2N)!}{2^N N!}$$

Other Questions

- How many ways can re-order the dancing partners so that k pairs are not found in the original pairing?

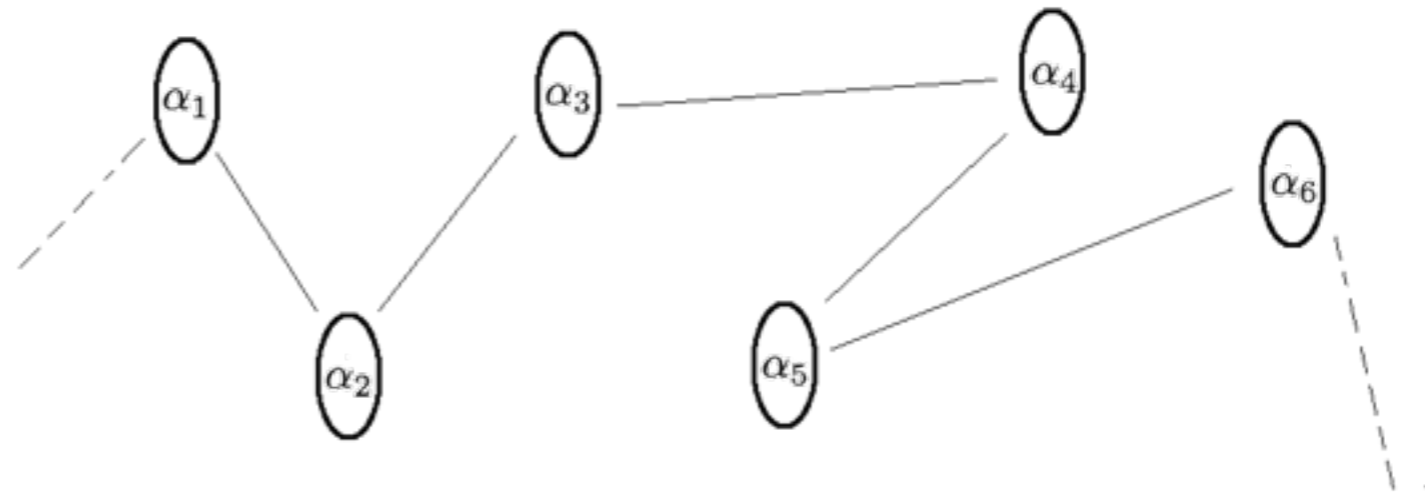
$$A_k = (-1)^k \sum_{\ell=0}^k (-1)^\ell \binom{k}{\ell} \prod_{i=1}^{\ell} (2i - 1)$$

where $\sum_{k=0}^N A_k = \frac{(2N)!}{2^N N!}$

Main Point
 There is an existing mathematical formalism to answer questions about rearranging an initial collection of pairs

Favorable Contact Model: Microstates

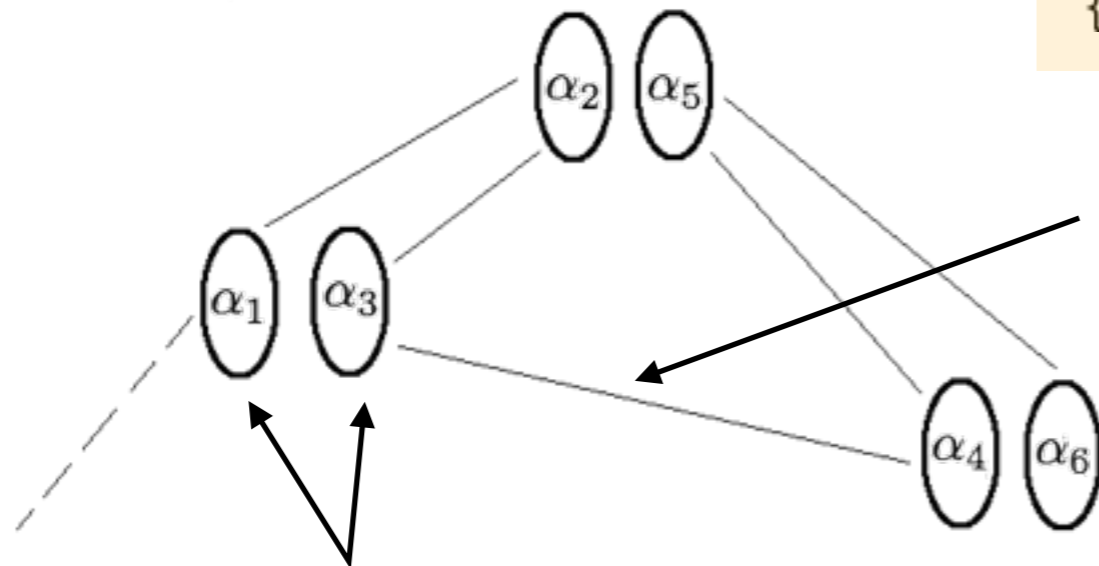
“Unfolded” configuration of 6 regions



Example of a microstate, i.e., a folded configuration of 6 regions

– regions are labeled as α_i with $i = 1, \dots, 2N$

– there is a single zero-energy (i.e., native) collection of pairwise contacts given by—for $2N$ regions— $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$.



– links between regions are infinitely extendible

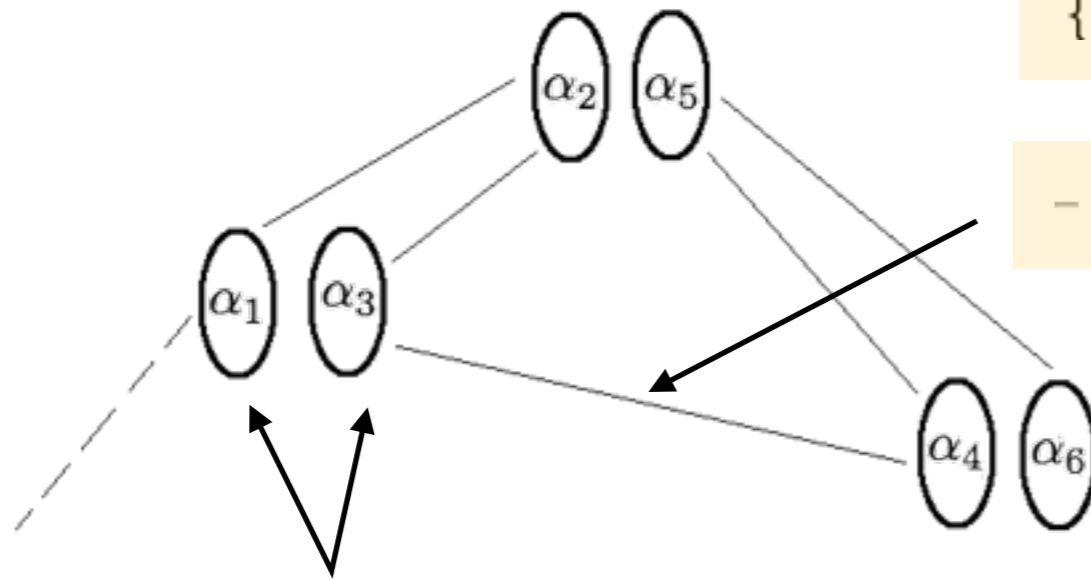
– only pairwise contacts between regions are possible

– microstates of system are given by collections of pairwise contacts (e.g., this microstate is $\{(\alpha_1, \alpha_3), (\alpha_2, \alpha_5), (\alpha_4, \alpha_6)\}$)

Favorable Contact Model: Energy

Example of a folded configuration of 6 regions

– regions are labeled as α_i with $i = 1, \dots, 2N$



– there is a single zero-energy (i.e., native) collection of pairwise contacts given by—for $2N$ regions— $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$.

– links between regions are infinitely extendible

– only pairwise contacts between regions are possible

– microstates of system are given by collections of pairwise contacts (e.g., this microstate is $\{(\alpha_1, \alpha_3), (\alpha_2, \alpha_5), (\alpha_4, \alpha_6)\}$)

Energy Definition:

– The microstate $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$ has zero energy. For an arbitrary microstate, there is an energy cost of γ_i if the contact pair i is not an element of the set of pairs $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$.

Example: 6 components

| State | Energy |
|--|----------------------------------|
| $(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), (\alpha_5, \alpha_6)$ | 0 |
| $(\alpha_1, \alpha_2), (\alpha_5, \alpha_4), (\alpha_3, \alpha_6)$ | $\gamma_2 + \gamma_3$ |
| $(\alpha_3, \alpha_6), (\alpha_1, \alpha_2), (\alpha_4, \alpha_5)$ | $\gamma_1 + \gamma_3$ |
| $(\alpha_1, \alpha_5), (\alpha_2, \alpha_4), (\alpha_3, \alpha_6)$ | $\gamma_1 + \gamma_2 + \gamma_3$ |
| $(\alpha_2, \alpha_5), (\alpha_3, \alpha_1), (\alpha_4, \alpha_6)$ | $\gamma_1 + \gamma_2 + \gamma_3$ |

Favorable Contact Model: Statistical Mechanics

System Assumptions:

- regions are labeled as α_i with $i = 1, \dots, 2N$
- only pairwise contacts between regions are possible
- links between regions are infinitely extendible
- microstates of system are given by collections of pairwise contacts (e.g., for $N = 3$, a microstate is $\{(\alpha_1, \alpha_3), (\alpha_2, \alpha_5), (\alpha_4, \alpha_6)\}$)
- there is a single zero-energy (i.e., native) collection of pairwise contacts given by—for $2N$ regions— $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$.

Energy Definition:

- The microstate $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$ has zero energy. For an arbitrary microstate, there is an energy cost of γ_i if the contact pair i is not an element of this set of pairs.

$$\mathcal{H}_N(\{(\theta_1^{(i)}, \theta_2^{(i)})\}) = \sum_{i=1}^N \gamma_i I(\theta_1^{(i)}, \theta_2^{(i)})$$

where

$$I(\theta_1^{(i)}, \theta_2^{(i)}) = \begin{cases} 0 & \text{if } (\theta_1^{(i)}, \theta_2^{(i)}) \in \{(\alpha_1, \alpha_2), \dots, (\alpha_{2N-1}, \alpha_{2N})\} \\ 1 & \text{otherwise.} \end{cases}$$

Example: 6 components

$\langle I(\theta_1^{(i)}, \theta_2^{(i)}) \rangle$: average number of non-native contacts

| State | # of non-native contacts | Energy |
|--|--------------------------|----------------------------------|
| $(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), (\alpha_5, \alpha_6)$ | 0 | 0 |
| $(\alpha_1, \alpha_2), (\alpha_5, \alpha_4), (\alpha_3, \alpha_6)$ | 2 | $\gamma_2 + \gamma_3$ |
| $(\alpha_3, \alpha_6), (\alpha_1, \alpha_2), (\alpha_4, \alpha_5)$ | 2 | $\gamma_1 + \gamma_3$ |
| $(\alpha_1, \alpha_5), (\alpha_2, \alpha_4), (\alpha_3, \alpha_6)$ | 3 | $\gamma_1 + \gamma_2 + \gamma_3$ |
| $(\alpha_2, \alpha_5), (\alpha_3, \alpha_1), (\alpha_4, \alpha_6)$ | 3 | $\gamma_1 + \gamma_2 + \gamma_3$ |

What are the equilibrium statistical mechanical properties of this system at an arbitrary T ?

Favorable Contact Model: Statistical Mechanics

Energy Definition:

- The microstate $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$ has zero energy. For an arbitrary microstate, there is an energy cost of γ_i if the contact pair i is not an element of this set of pairs.

$$\mathcal{H}_N(\{(\theta_1^{(i)}, \theta_2^{(i)})\}) = \sum_{i=1}^N \gamma_i I(\theta_1^{(i)}, \theta_2^{(i)})$$

where

$$I(\theta_1^{(i)}, \theta_2^{(i)}) = \begin{cases} 0 & \text{if } (\theta_1^{(i)}, \theta_2^{(i)}) \in \{(\alpha_1, \alpha_2), \dots, (\alpha_{2N-1}, \alpha_{2N})\} \\ 1 & \text{otherwise.} \end{cases}$$

What are the equilibrium statistical mechanical properties of this system at an arbitrary T ?



What is the partition function for this system?

$$Z_N(\{\beta\gamma_i\}) = \sum_{\{(\theta_1^{(i)}, \theta_2^{(i)})\}} \exp\left(-\beta \sum_{j=1}^N \gamma_j I(\theta_1^{(j)}, \theta_2^{(j)})\right)$$

(Some mathematical work, i.e., dancing partners)

Allows us to compute:

- two-region and four-region correlations
- free energy as a function of deviation from native state
- native-to-unfolded transition temperatures

$$Z_N(\{\beta\gamma_i\}) = \frac{1}{\sqrt{\pi}} \int_0^\infty dt t^{-1/2} e^{-t} \prod_{k=1}^N \left(1 + (2t - 1)e^{-\beta\gamma_k}\right)$$

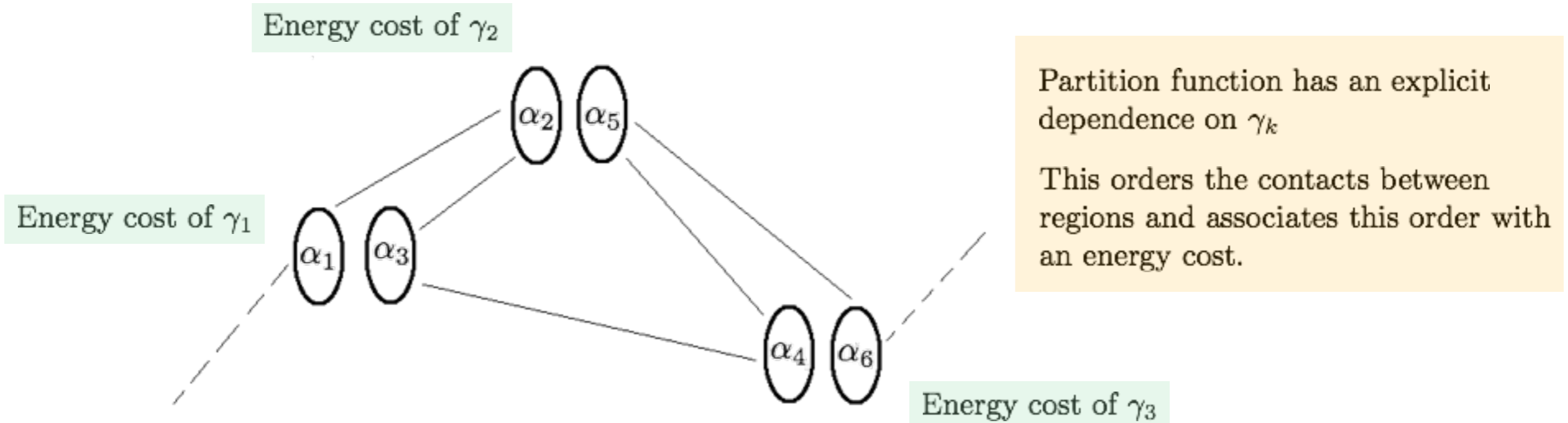
An unphysical but analytically soluble folding model

* Very similar in form to the result from the statistical mechanics of the symmetric group:

$$Z_N(\{\beta\lambda_i\}) = \int_0^\infty ds e^{-s} \prod_{\ell=1}^N \left[1 + (s - 1)e^{-\beta\lambda_\ell}\right]$$

FCM: Unphysical Assumption

Favorable Contact Partition Function:
$$Z_N(\{\beta\gamma_i\}) = \frac{1}{\sqrt{\pi}} \int_0^\infty dt t^{-1/2} e^{-t} \prod_{k=1}^N \left(1 + (2t - 1)e^{-\beta\gamma_k}\right)$$



Partition function has an explicit dependence on γ_k

This orders the contacts between regions and associates this order with an energy cost.

...but the energy cost should depend on **the properties** of the adjacent regions and *not* on the arbitrary ordering of the contacts.

For solubility, our model does not consider the **specific properties** of interactions between regions...

...but we can consider **distributions** of these interaction properties.

→ Each γ_k is drawn from distribution $\rho_0(\gamma)$

FCM: Parameterizing Ignorance

Favorable Contact Partition Function:
$$Z_N(\{\beta\gamma_i\}) = \frac{1}{\sqrt{\pi}} \int_0^\infty dt t^{-1/2} e^{-t} \prod_{k=1}^N \left(1 + (2t - 1)e^{-\beta\gamma_k}\right)$$

Instead of giving each contact k
an energy cost of γ_k

we say the energy cost for contact k
is drawn from a distribution $\rho_0(\gamma)$

Example: Gaussian Distribution of Energy Costs

Each γ_k is drawn from a normal distribution
with mean γ_0 and variance σ_γ^2 .

$$\gamma_k \sim \mathcal{N}(\gamma_0, \sigma_\gamma^2)$$

Mean energy cost; proxy for
the **stability** of the native
configuration

If $\gamma_0 \uparrow$, then stability increases
If $\gamma_0 \downarrow$, then stability decreases

Variance of energy cost;
proxy for the **heterogeneity**
of the chain of regions

If $\sigma_\gamma \uparrow$, then chain is more heterogeneous
If $\sigma_\gamma \downarrow$, then chain is more homogeneous

Question

How do we determine the **average properties**
of a system defined by such a distribution of
interaction energies?

Answer

Compute the quenched
free energy!

Generalized Favorable Contact Model

Favorable Contact Free Energy
(with energy cost distributions)



Quenched Free Energy

$$\langle \ln Z_N(\{\beta\gamma_i\}) \rangle = \int_{-\infty}^{\infty} \prod_{j=1}^N d\gamma_j \rho_0(\gamma_j) \ln \frac{1}{\sqrt{\pi}} \int_0^{\infty} dt t^{-1/2} e^{-t} \prod_{k=1}^N \left(1 + (2t-1)e^{-\beta\gamma_k} \right)$$

“Generalized Favorable Contact Model”

OK. It’s a complicated mathematical expression. So what?

This favorable contact free energy is similar in form to that for the **permutation glass**

$$-\beta F_{\text{perm. glass}} = \int_{-\infty}^{\infty} \prod_{j=1}^N d\lambda_j \rho_0(\lambda_j) \ln \int_0^{\infty} ds e^{-s} \prod_{k=1}^N \left(1 + (s-1)e^{-\beta\lambda_k} \right)$$

Permutation Glass (definition)
physical system where the state space consists of permutations of a list and the Hamiltonian depends on random parameters. (similar to spin glass)

Physical Results of Permutation Glass

For a chain length N , a non-degenerate free energy minimum exists only if

$$\frac{\langle \lambda \rangle}{\sigma_\lambda} \gtrsim \sqrt{2 \ln N}$$

Can we establish similar constraints for the generalized favorable contact model?

mean/variance of energy cost for deviating from non-degenerate microstate

e.g.

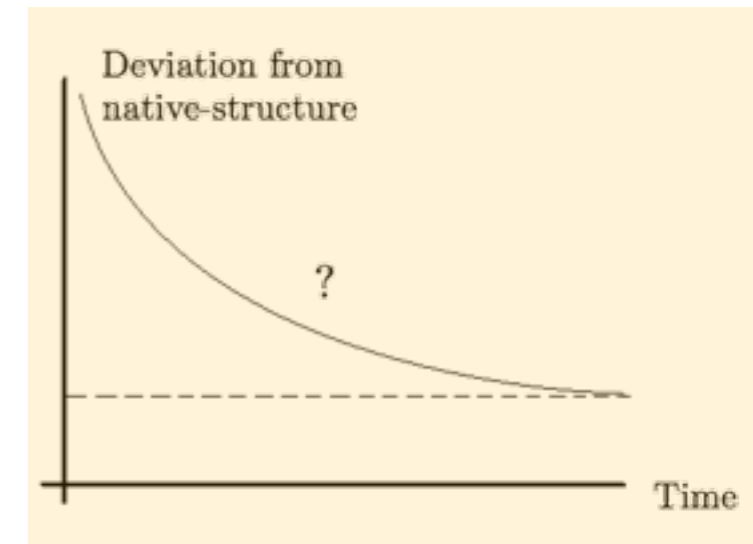
| State | Energy |
|----------------------------------|-------------------------------------|
| $(\omega_1, \omega_2, \omega_3)$ | 0 |
| $(\omega_2, \omega_1, \omega_3)$ | $\lambda_1 + \lambda_2$ |
| $(\omega_3, \omega_2, \omega_1)$ | $\lambda_1 + \lambda_3$ |
| $(\omega_1, \omega_3, \omega_2)$ | $\lambda_2 + \lambda_3$ |
| $(\omega_2, \omega_3, \omega_1)$ | $\lambda_1 + \lambda_2 + \lambda_3$ |
| $(\omega_3, \omega_1, \omega_2)$ | $\lambda_1 + \lambda_2 + \lambda_3$ |

where each λ_i is drawn from a probability distribution

Generalized FCM: Questions

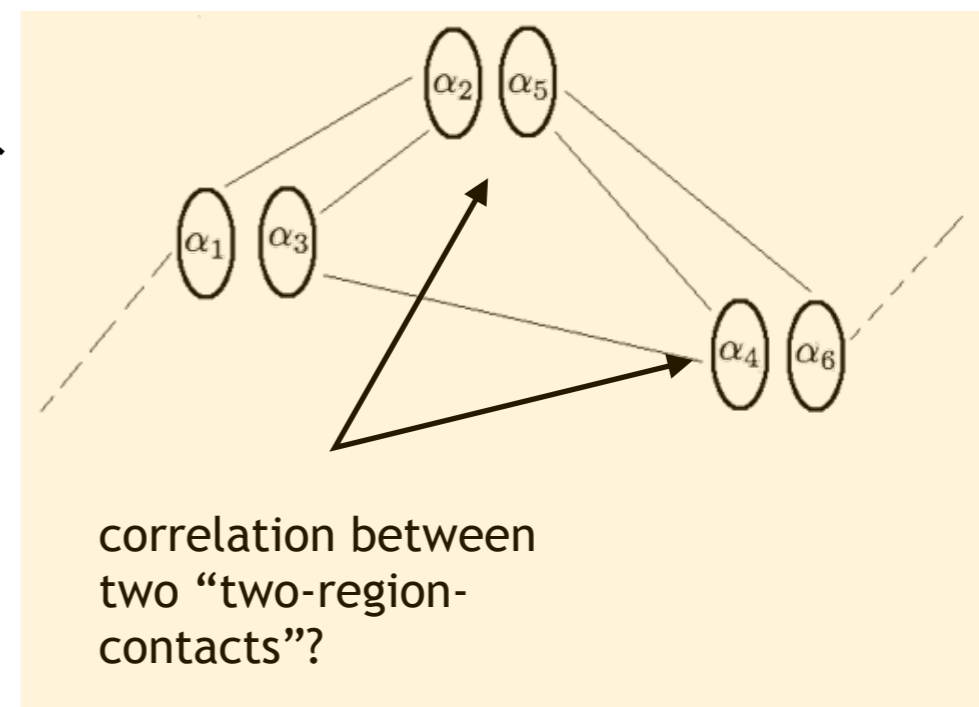
Questions to ask about General Favorable Contact Model

- **Existence of Native State:** How do the stability and heterogeneity properties of this abstracted polypeptide affect the existence of a native state?
- **Relaxation (“Folding”) Time:** How much time does it take the system to reach thermal equilibrium? Such a time can be taken to be how long it takes the abstracted polypeptide to “fold.”
- **Non-native contacts:** How does the average number of non-native contacts vary with temperature and parameters?
- **Regional Contact Correlations:** What is the four-region correlation? (i.e., the correlation between two different two-region contact regions)



$\langle I(\theta_1^{(i)}, \theta_2^{(i)}) \rangle$: average number of non-native contacts

$\langle \gamma \rangle$: proxy for stability of native state
 σ_γ : proxy for heterogeneity of chain
 \implies How does this affect the existence of native state?



Generalized FCM: Unphysical Assumptions

Unphysical assumptions of model

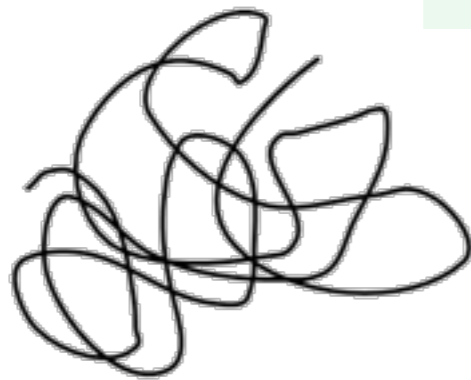
- **Even number of regions:** Assumed the chain consisted of exactly an even number of regions.
- **No specific interaction matrix:** In lieu of an interaction matrix, which could account for how specific regions interact with each other, we assumed the interaction energies could be modeled by quenched disorder.
- **Infinitely extendable chain:** To allow for all possible interacting combinations, we allowed our abstracted polypeptide to be infinite extendable.
- **Only includes interacting pairs:** Although two regions of a polypeptide chain can sometimes define interaction, they do not exclusively do so.

Connecting Protein Folding and Protein Design

(Abstracted) Design and Folding Problems

Study a model of possible sequence orders for a given presumed structure

$$Z_N^{\text{perm.}}(\{\beta\lambda_i\}) = \int_0^\infty ds e^{-s} \prod_{k=1}^N (1 + (s-1)e^{-\beta\lambda_k})$$



Structure

Protein Design



... A - R - H - G - L - H ...



Sequence

Protein Folding

Study a model of possible pairwise contacts for a given sequence of contact regions

$$Z_N^{\text{FCM}}(\{\beta\gamma_i\}) = \frac{1}{\sqrt{\pi}} \int_0^\infty dt t^{-1/2} e^{-t} \prod_{k=1}^N (1 + (2t-1)e^{-\beta\gamma_k})$$



END