

# Experimental Realizations of Favorable-Contact Model

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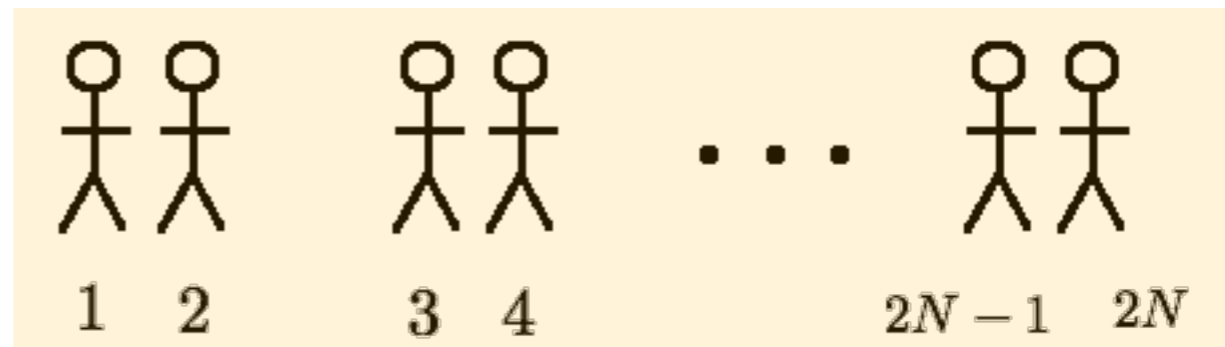
*Mobolaji Williams — Shakhnovich Group Meeting— Aug. 15, 2017*

*(Thanks to Evgeny Serebryany)*

# Recall: “Dancing Partners Problem”

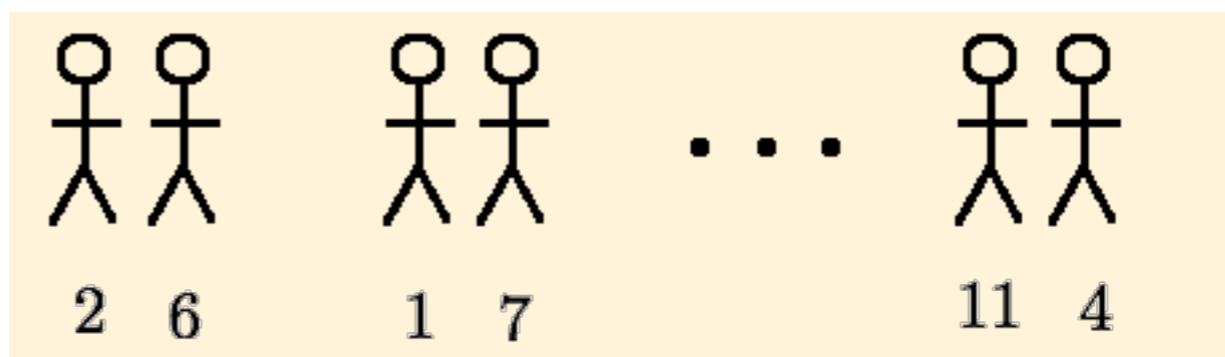
(B. H. Margolius, “Avoiding your spouse at a bridge party, 2001)

- ▶  $N$  partners (i.e.,  $2N$  total people) arrive at a dancing party



↓ (reorder)

- ▶ Each dance partner pair is re-ordered such that each person may or may not be with their original dance partner.



- ▶ In how many ways can this happen?  
(What is the size of the state space?)

$$|\mathcal{S}_N| = \frac{(2N)!}{2^N N!}$$

**Main Question:** Can we develop an experimentally testable physical model from these statistics?

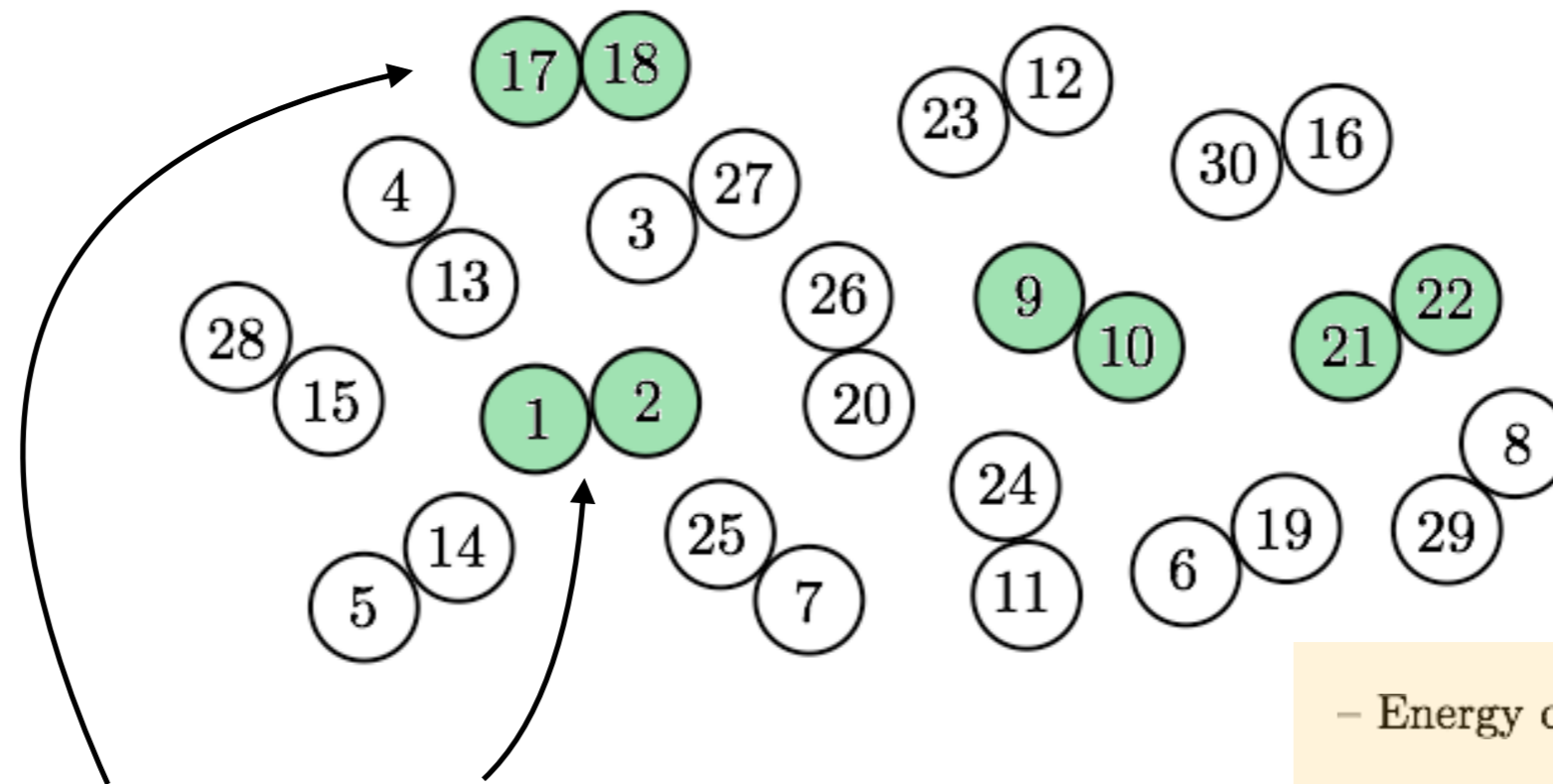
# Basics of Favorable Contact Model

Say we have a system of  $2N$  distinguishable elements which are organized into pairs.

We denote the elements as  $\alpha_1, \alpha_2, \alpha_3, \dots, \alpha_{2N}$ , and each element can be paired with any other element and the ordering within the pair is not important.

## Example Microstate:

– A possible microstate of the  $2N = 30$  system



– Energy of microstate is  $E = 11\gamma_0$

We also take there to be a single zero-energy collection of pairings. We call these pairings "favorable contacts", and we write this zero-energy microstate as  $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$ .

For other microstates, whenever a pairing of elements does not pair  $\alpha_{2k-1}$  with  $\alpha_{2k}$ , (i.e., for each unfavorable contact) there is an energy cost of  $\gamma_0$ .

# Equilibrium Thermodynamics of Model

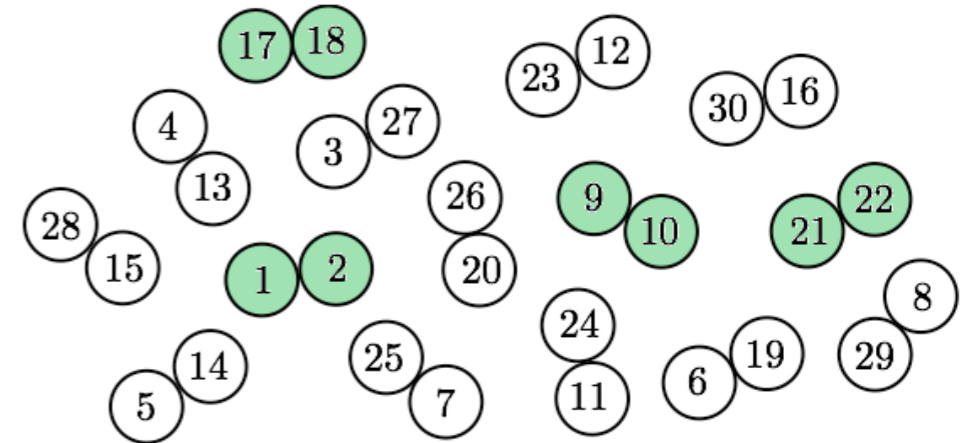
The partition function for this system is then

$$Z_N(\beta\gamma_0) = \sum_{\{(\theta_1^{(i)}, \theta_2^{(i)})\}} \exp \left( -\beta\gamma_0 \sum_{j=1}^N I(\theta_1^{(j)}, \theta_2^{(j)}) \right),$$

where  $(\theta_1^{(i)}, \theta_2^{(i)})$  denotes the 1st and 2nd components (where the internal ordering of the components is not important) of the  $i$ th interaction pair, and where

$$I(\theta_1^{(j)}, \theta_2^{(j)}) = \begin{cases} 0 & (\theta_1^{(j)}, \theta_2^{(j)}) = (\alpha_1, \alpha_2), \dots, \text{ or } (\alpha_{2N-1}, \alpha_{2N}) \\ 1 & \text{otherwise.} \end{cases}$$

What are the equilibrium thermodynamics of this model?



Redefine partition function in terms of  $\ell$ , the number of unfavorable contacts.

$$Z_N(\beta\gamma_0) = \sum_{\ell=0}^N \binom{N}{\ell} a_\ell e^{-\beta\gamma_0 \ell}$$

combinatorics manipulations

$\ell$  : number of unfavorable contacts

(Definition of  $a_\ell$ ):  $2\ell$  objects are initially paired in some way and then are re-paired so that no object is with its original partner.  $a_\ell$  is the number of ways this re-pairing can occur.

$$Z_N(\beta\gamma_0) = \frac{1}{\sqrt{\pi}} \int_0^\infty dt t^{-1/2} e^{-t} \left[ 1 + (2t - 1)e^{-\beta\gamma_0} \right]^N$$

$$a_\ell = (-1)^\ell \sum_{k=0}^{\ell} (-1)^k \binom{\ell}{k} \prod_{i=1}^k (2i - 1)$$

# Physics Predictions: Order Parameter and $T_c$

$$Z_N(\beta\gamma_0) = \frac{1}{\sqrt{\pi}} \int_0^\infty dt t^{-1/2} e^{-t} \left[ 1 + (2t - 1)e^{-\beta\gamma_0} \right]^N$$

What physical predictions  
can we obtain from the model?

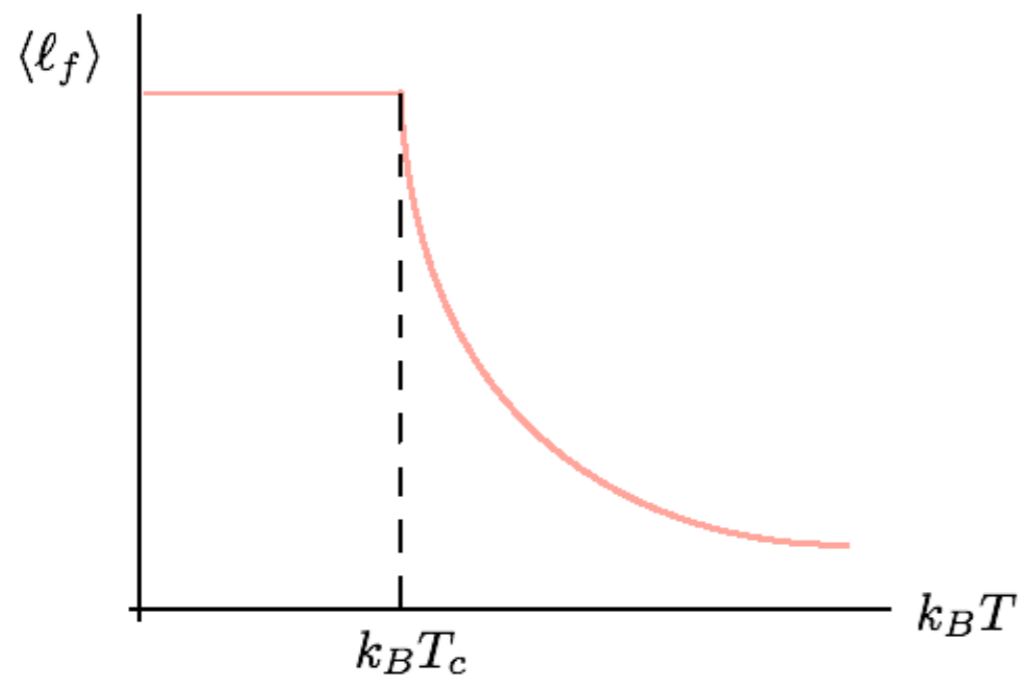
$$\langle \ell_f \rangle = N + \frac{\partial}{\partial(\beta\gamma_0)} \ln Z_N(\beta\gamma_0)$$

$$\simeq \frac{1}{4(1 - e^{-\beta\gamma_0})} \left[ 2(N - 1) + e^{\beta\gamma_0} - \sqrt{(2N - e^{\beta\gamma_0})^2 - 4(e^{\beta\gamma_0} - 1)} \right]$$

**Order Parameter:**  $\langle \ell \rangle$ , average number of unfavorable contacts

\*  $\langle \ell_f \rangle = N - \langle \ell \rangle$ , average number of favorable contacts

Schematic plot of average number of favorable contacts



Temperature where  $\langle \ell_f \rangle \in \mathbb{R} \rightarrow \langle \ell_f \rangle \in \mathbb{C}$ :

$$k_B T_c = \frac{\gamma_0}{\ln(2N)} \left( 1 + \sqrt{\frac{2}{N \ln(2N)}} \right) + \mathcal{O}(N^{-1})$$

Checking Limiting cases

- $\gamma_0 \rightarrow \infty, T_c \rightarrow \infty$ : Favorable contacts persist at higher temperature
- $N \rightarrow \infty, T_c \rightarrow 0$ : Larger entropic contributions require a lower  $T_c$

We can compute this critical temperature.

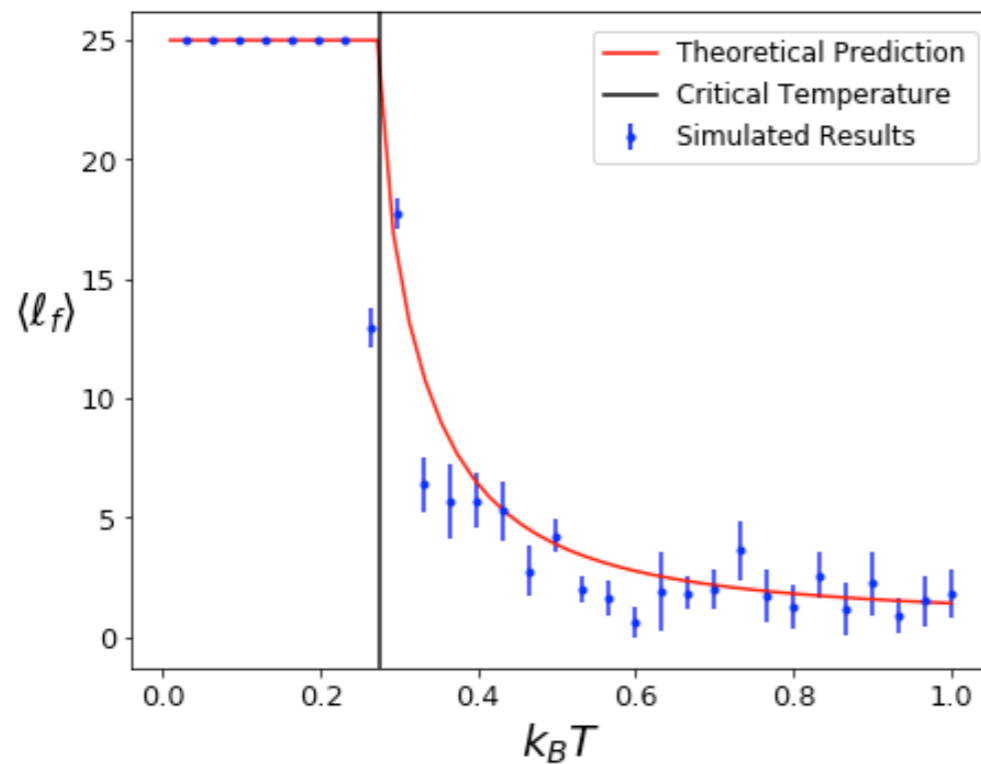
# Comparison with Simulations

$$\langle \ell_f \rangle \simeq \begin{cases} N & \text{for } T < T_c, \\ \frac{1}{4(1 - e^{-\beta\gamma_0})} \left[ 2(N - 1) + e^{\beta\gamma_0} - \sqrt{(2N - e^{\beta\gamma_0})^2 - 4(e^{\beta\gamma_0} - 1)} \right] & \text{for } T \geq T_c. \end{cases}$$

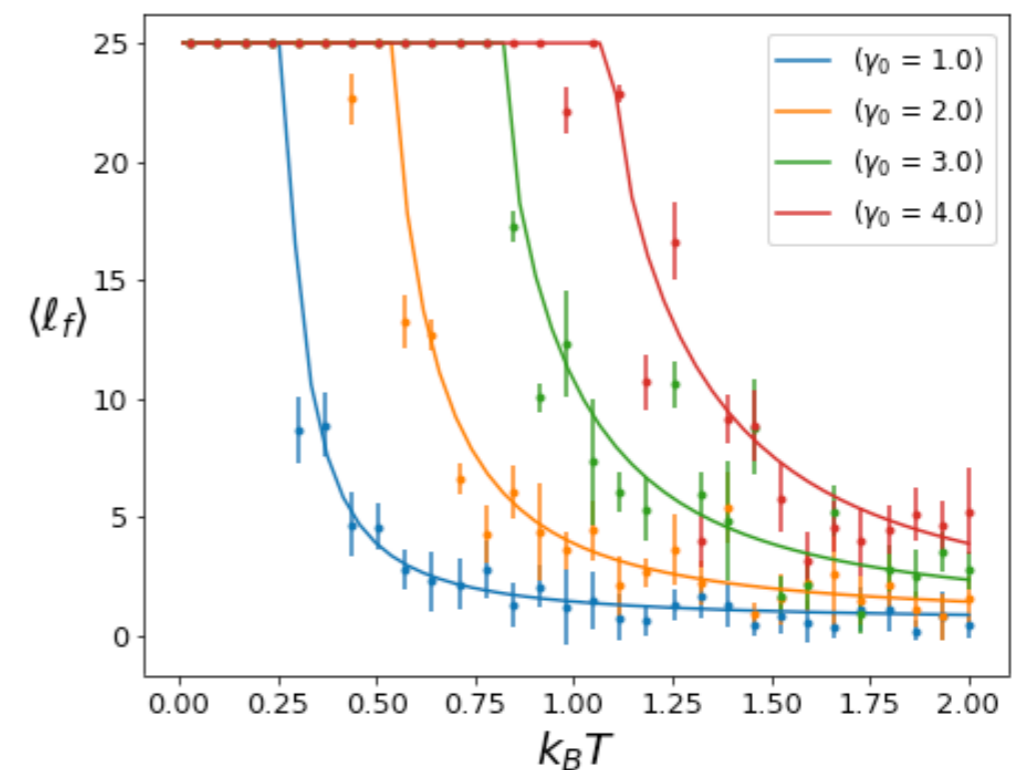
( $\langle \ell_f \rangle = N - \langle \ell \rangle$ , average number of favorable contacts)

$$k_B T_c = \frac{\gamma_0}{\ln(2N)} \left( 1 + \sqrt{\frac{2}{N \ln(2N)}} \right) + \mathcal{O}(N^{-1})$$

(critical temperature; i.e.,  $T$  where  $\langle \ell_f \rangle \in \mathbb{R} \rightarrow \langle \ell_f \rangle \in \mathbb{C}$ )



Monte Carlo Simulation for  $N = 25$  pairs,  $\gamma_0 = 1.0$



Monte Carlo Simulation for  $N = 25$  pairs, for various  $\gamma$

# Physics Predictions: Heat Capacity and Critical Behavior

Does the model exhibit critical behavior?

(Definition of universality):  
 When the behavior of a system near a critical transition is independent of the microscopic details of the system.

– critical exponents indicate the type of universality a system exhibits

## Determining critical exponents

– How do order parameter and heat capacity vary with  $T$  near critical temperature (i.e., for  $|T - T_c| \ll T_c$ )?

$$\langle \ell_f \rangle \sim N - (T - T_c)^\beta, \quad C_T \sim (T - T_c)^{-\alpha}.$$

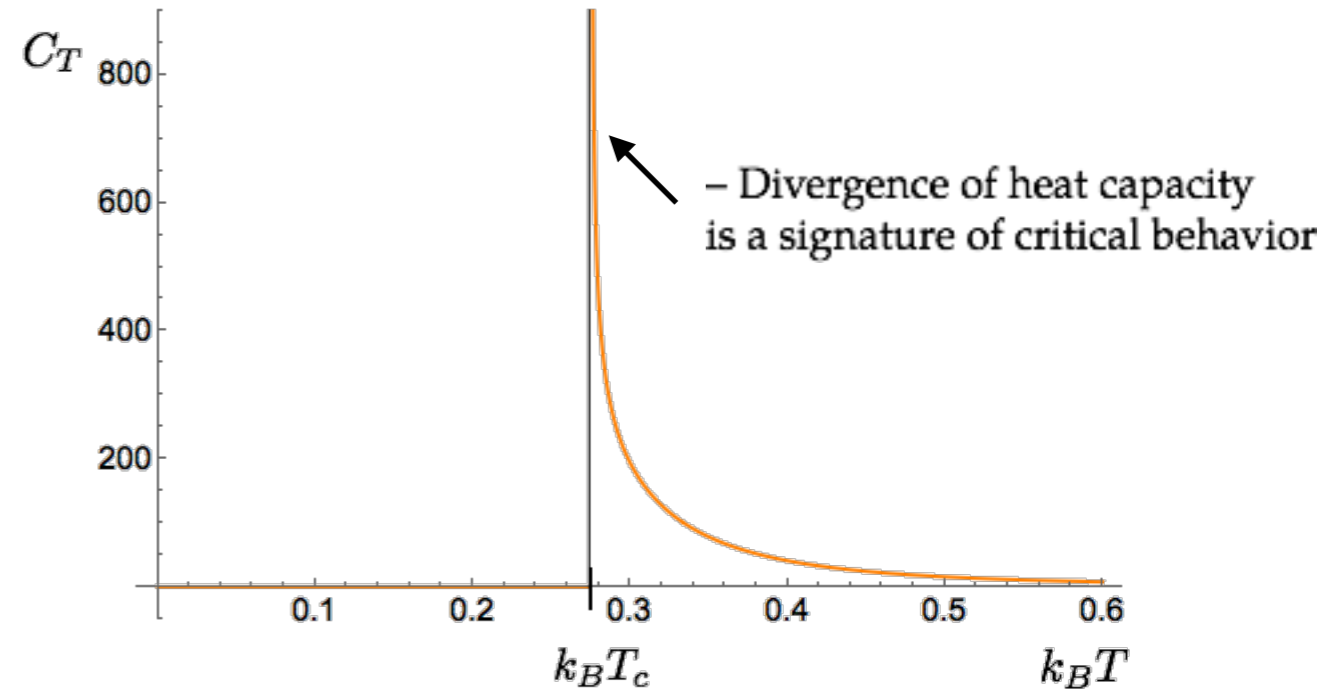
$$\rightarrow \langle \ell_f \rangle \simeq N - \frac{1}{2} (2N)^{1/4} |f'(T_c)|^{1/2} (T - T_c)^{1/2}$$

$$C_T \simeq \frac{\gamma_0}{4} (2N)^{1/4} |f'(T_c)|^{1/2} (T - T_c)^{-1/2}$$

$$* f(T) = e^{\gamma_0/k_B T}$$

$$\alpha = \beta = 1/2$$

Heat Capacity for  $N = 25$  and  $\gamma_0 = 1.0$



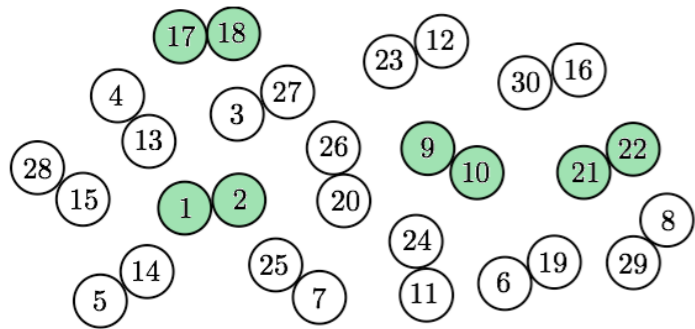
– Critical behavior is important because of what it suggests about universality

Table of critical exponents for canonical systems

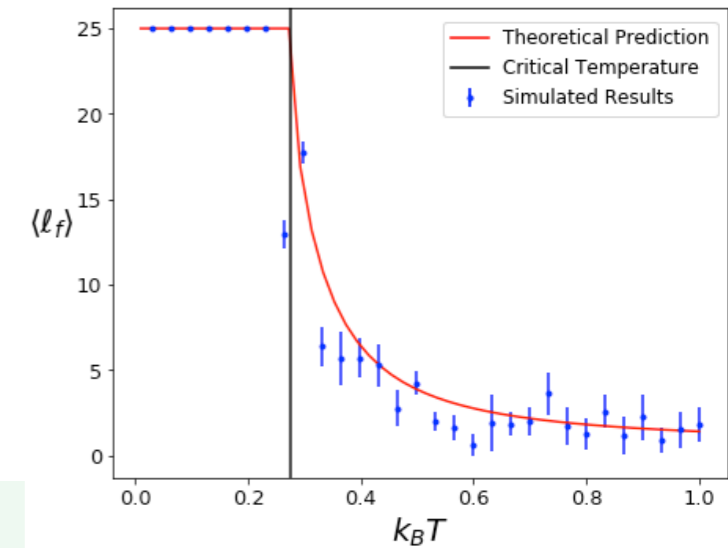
Theory	$\alpha$	$\beta$
Mean Field Theory Ising	0	1/2
Liquid Gas Transition	0	1/2
2D Ising Model	0	1/8
Favorable contact Model	1/2	1/2

Critical behavior, but no existing universality class!

# Physics Predictions: Summary



Model has a number of physical predictions (which can be checked by simulations)



## 1. Order Parameter:

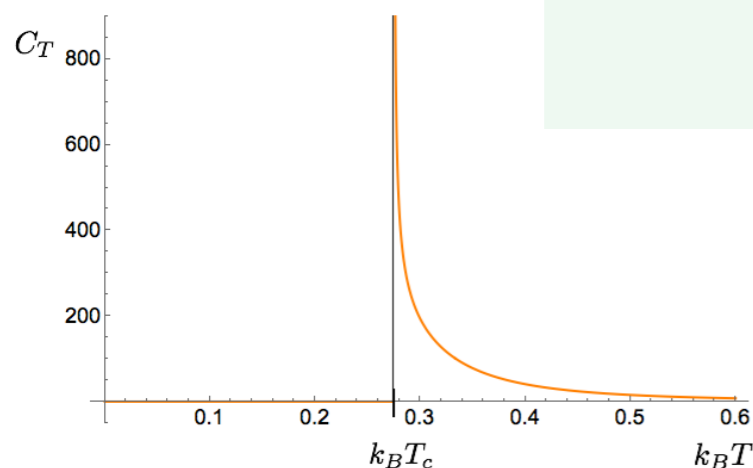
$$\langle \ell_f \rangle \simeq \frac{1}{4(1 - e^{-\beta\gamma_0})} \left[ 2(N - 1) + e^{\beta\gamma_0} - \sqrt{(2N - e^{\beta\gamma_0})^2 - 4(e^{\beta\gamma_0} - 1)} \right].$$

## 2. Critical Temperature:

$$k_B T_c = \frac{\gamma_0}{\ln(2N)} \left( 1 + \sqrt{\frac{2}{N} \frac{1}{\ln(2N)}} \right) + \mathcal{O}(N^{-1}).$$

## 3. Critical Exponents:

$$C_T \sim (T - T_c)^{-1/2} \text{ and } \langle \ell \rangle \sim (T - T_c)^{1/2}.$$



Are there any experimental systems where we can test these ideas?



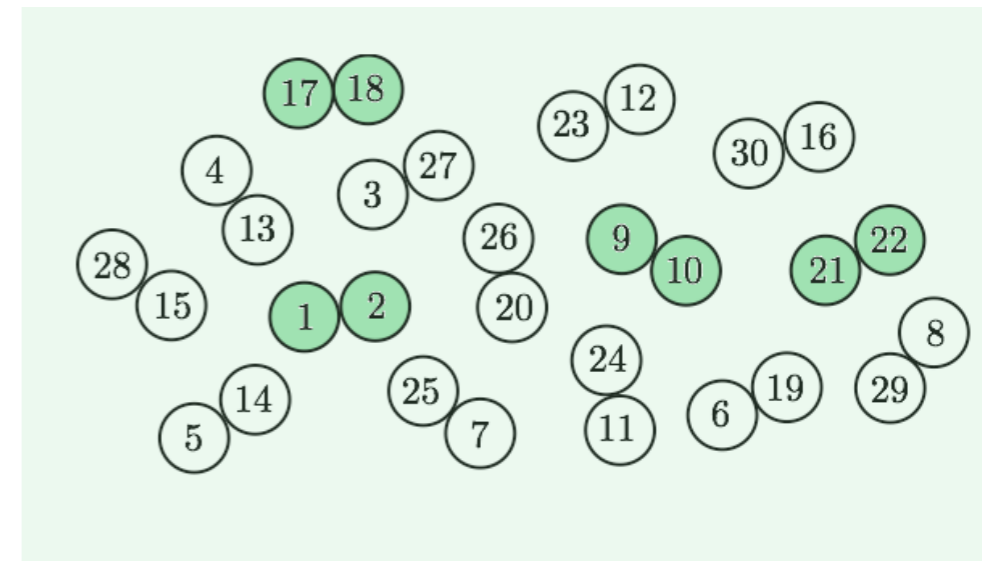
# Needed Elements

Is it possible to prepare a system where:

1. there is an even number of distinguishable subunits
2. subunits exist primarily subunit-subunit pairs (energetically unfavorable to exist otherwise)
3. there is a single optimal collection of pairings
4. there is a measurable energy cost for deviations from these optimal pairings.

$2N$  subunits

$$\mathcal{H} = \gamma_0 \sum_{j=1}^N I(\theta_1^{(j)}, \theta_2^{(j)})$$



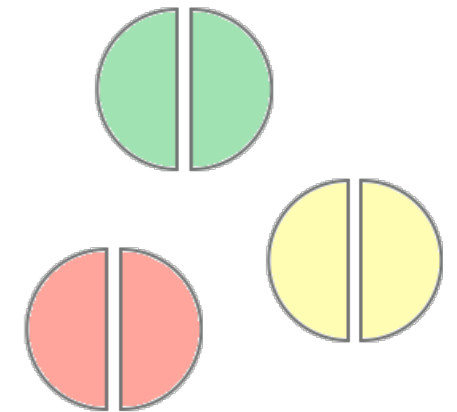
**Optimal Pairings:**  $\{(\alpha_1, \alpha_2), (\alpha_3, \alpha_4), \dots, (\alpha_{2N-1}, \alpha_{2N})\}$

# Candidate Biophysical Systems

## Possible Systems

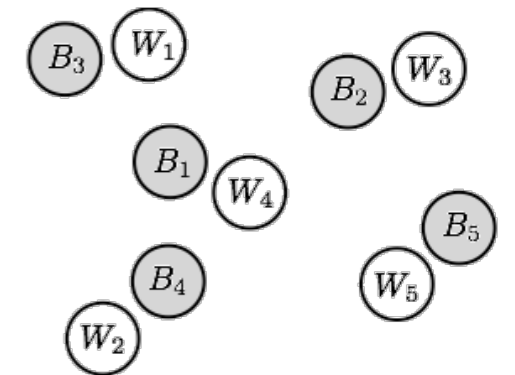
### 1. Protein-protein interactions

1. need to engineer/collect proteins prone to forming dimers
2. avoid formation of oligomers greater than 2
3. need to determine interaction energies for pairs



### 2. Transcription Factor-single stranded DNA (ssDNA) interactions

1. **Advantage:** ability to tune sequence allows for precise binding
2. **Problem:** This is a “gendered” favorable contact model
  - Analyze using previous permutation model



### 3. ssDNA-ssDNA interactions

1. Don't have to deal with complications of tertiary structure
2. sequences provide simple biophysical estimate of model parameters
3. (supposedly) easy to order/prepare

e.g., Unfavorable contact

5'-AGCTAACGTA-3'  
3'-GGACTACTTT-5'

e.g., Favorable Contact

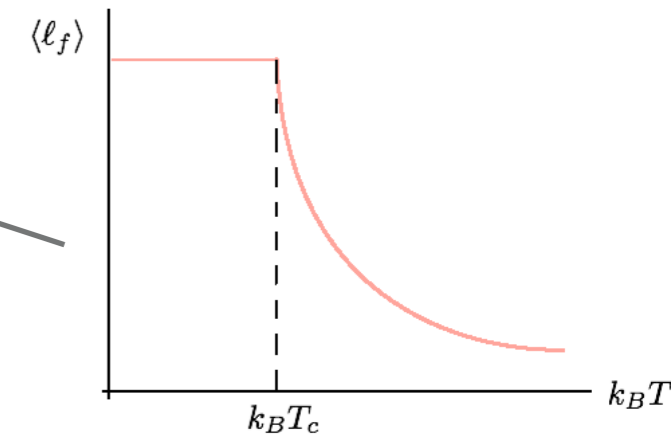
5'-AGCTAACGTA-3'  
3'-TCGATTGCAT-5'



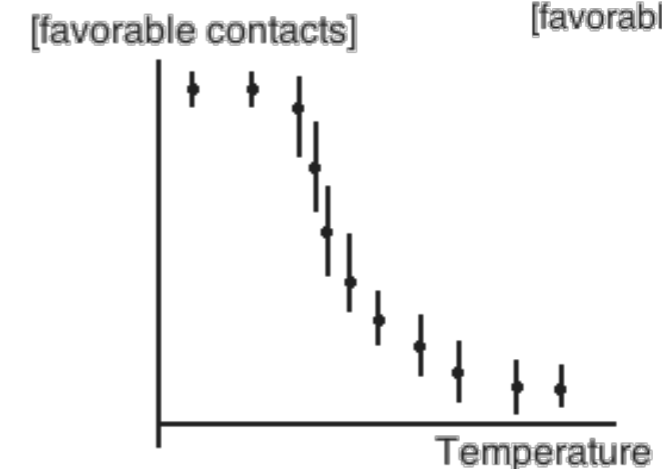
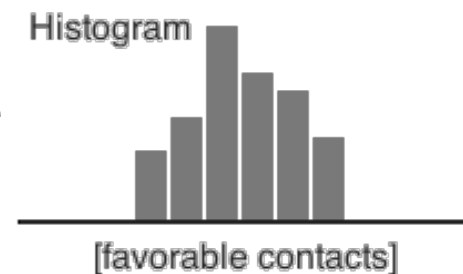
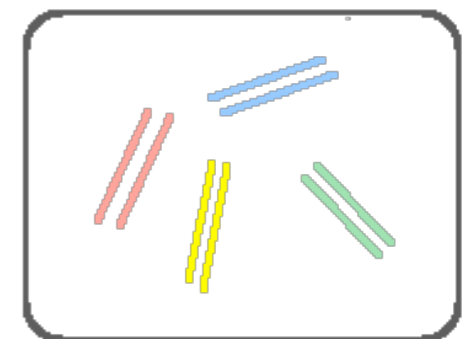
# Setup with Single Stranded DNA (ssDNA)

**Experimental Goal:** Measure number of ssDNA bonded to their complementary strands across a range of temperatures.

1. Prepare  $N$  different ssDNA  $n$ -mers
  1. Not clear how long they should be, but the shorter the better. (in order to prevent partial alignment of sequences, and formation of hairpins)
2. Prepare  $N$  complementary ssDNA fragments (all in an aqueous solution)
3. Begin at high temperature such that ssDNA are unbound
4. Rapidly lower temperature to desired  $T_i$
5. Measure concentration or fraction of favorable contacts\*
6. Repeat 1–5 at specific  $T_i$  to accumulate statistics
7. Repeat 1–6 for range of temperatures  $\{T_1, T_2, T_3, \dots\}$



5'-AGCTAACGTA-3'



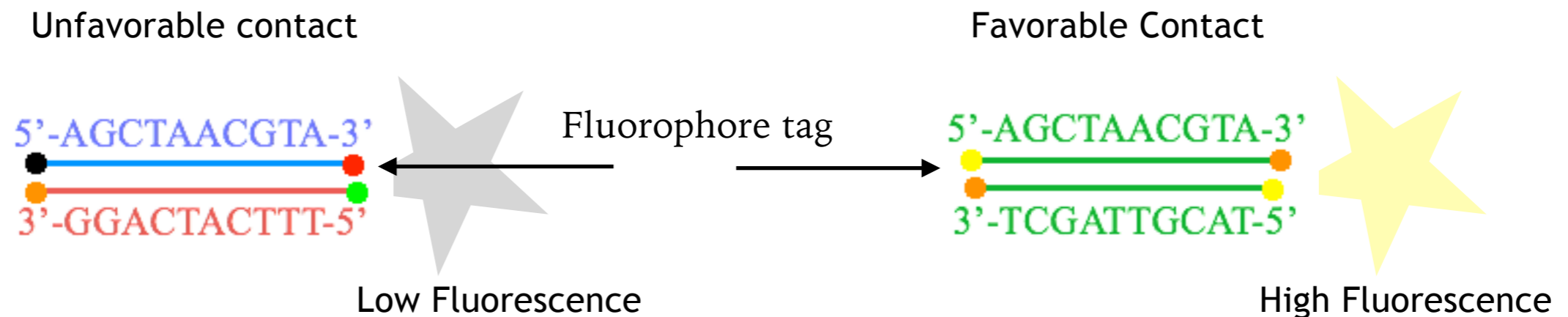
# Measuring Number of Favorable Contacts

**Pivotal Question:** How do we measure number of favorable contacts?

- Attaching fluorophores to ends of ssDNA can allow us to measure the amount of DNA that has formed double strands

(Quantifying DNA concentrations using fluorometry: A comparison of fluorophores, Rengarajan et. al. (2002) )

- Certain combinations of fluorophore tags lead to quenching or amplification of light produced by fluorophore



Fluorescence of Solution  $\sim$  Concentration of Favorable Contacts

# Short list of possible issues

$$Z_N^{\text{FCM}}(\{\beta\gamma_i\}) = \frac{1}{\sqrt{\pi}} \int_0^\infty dt t^{-1/2} e^{-t} \prod_{k=1}^N (1 + (2t - 1)e^{-\beta\gamma_k})$$

1. **Parameter Space:** The model only has two free parameters (energy cost and the number of subunits). Might not be enough freedom to represent an experimental situation.

Free parameters:  
 $\gamma_0$  and  $N$

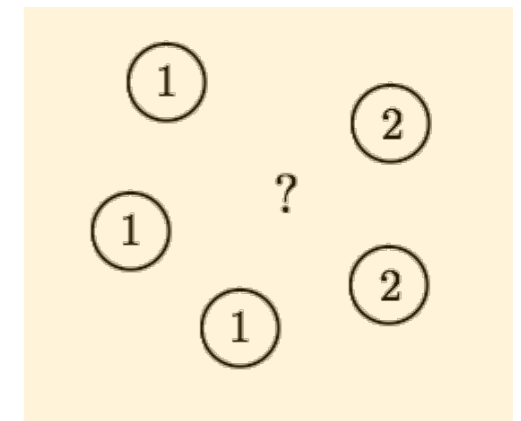
➤ Would have to consider the generalized Favorable Contact Model

2. **Fragment Overlap:** It's possible for complementary fragments to overlap in a way so that the nucleotide bases are not coupled to complementary base. Model doesn't allow for such partial contacts.



3. **Multiple Copies of Each ssDNA:** Experimentally, it is not feasible to make a single copy of each strand, but in the model there is only a single copy of each subunit.

➤ Could model be modified to include multiple subunits of each type?  
(Or can the model encompass this situation?)



4. **Biophysical Issues:** The pivotal physical prediction of the model is the critical temperature. Is this temperature in a physically relevant range?

$$k_B T_c = \frac{\gamma_0}{\ln(2N)} \left( 1 + \sqrt{\frac{2}{N \ln(2N)}} \right) + \mathcal{O}(N^{-1})$$

# “Order of Magnitude” Sanity Check

**Question:** Does an estimate of the critical temperature yield a reasonable temperature (e.g.,

1. below melting temperature of DNA and
2. above freezing temperature of water?)

$$k_B T_c \simeq \frac{\gamma_0}{\ln(2N)}$$

Energy cost for unfavorable contact

$N$ : number of favorable contact pairings

- Say we have a 30-mer oligonucleotide.
- For its double stranded form, we can approximate a melting temperature of:
- Let's approximate the energy cost as:
- Assume there are  $N = 25$  favorable pairings
- We then find a critical temperature estimate of

e.g., 5'-ACT GTG TCG TAT GAA TCG-3'

$$T_m \approx 60^\circ \text{C}$$

$$\gamma_0 = k_B T_m \approx 3 \text{ kJ/mol}$$

$$N = 25$$

$$T_c \simeq \frac{\gamma_0}{k_B \ln(2N)} \approx 92 \text{ K} = -181^\circ \text{C}$$

Rough prediction of critical temperature (in Kelvin) is **off by a factor of ~3** from a physical regime.



*END*